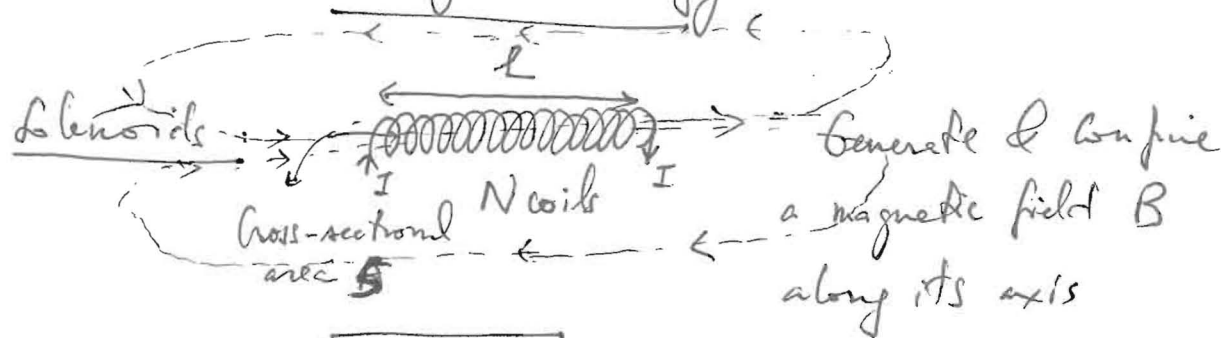


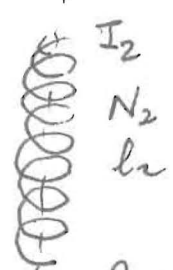
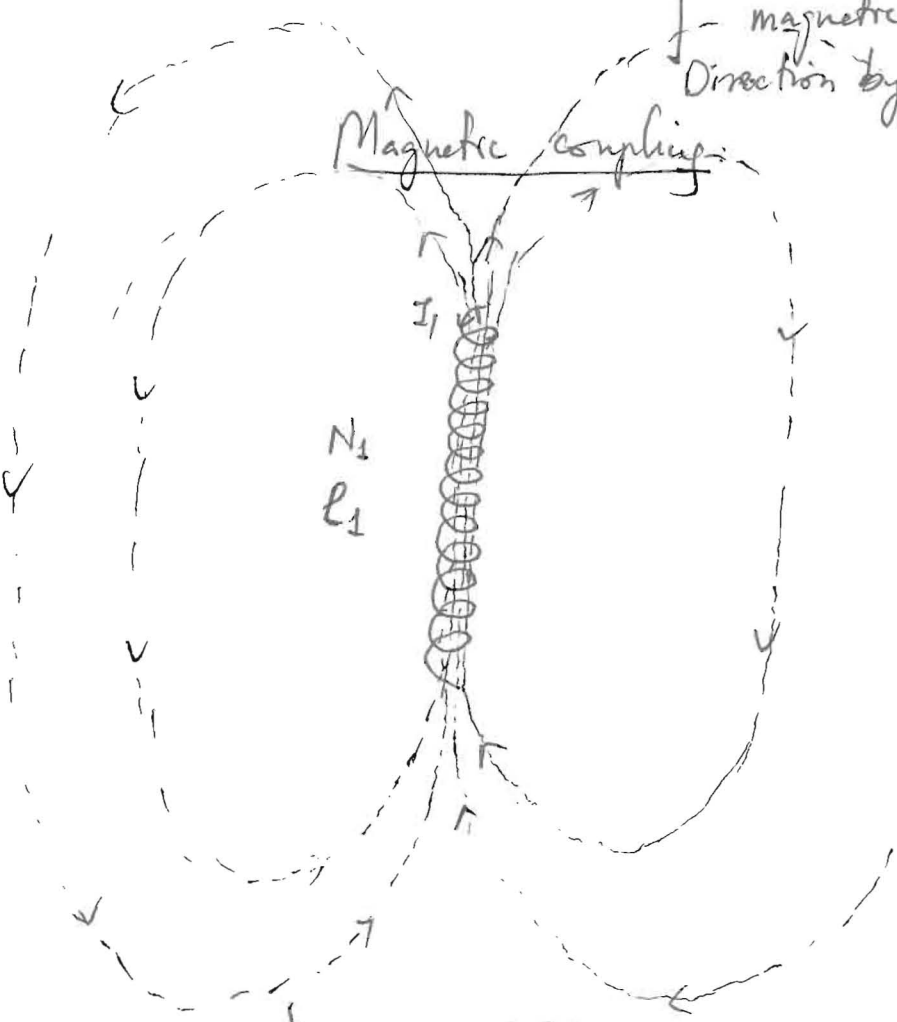
Self & Mutual Inductance (Magnetic Coupling) & Magnetic Energy



$$B = \mu n I$$

$n \equiv \frac{N}{l}$ or density of coils per unit length

magnetic permeability
Direction by RHR (RH fingers wrapping as current flows inside coil $\rightarrow$ thumb gives direction of $\vec{B}$)



$$B_2 = \mu n_2 I_2 ; n_2 \equiv \frac{N_2}{l_2}$$

(Field lines for B_2 are not shown)

- Magnetic field lines created by solenoid #1 affects solenoid #2. And vice versa!
 $\leftrightarrow$ mutual inductance.

M (unit in S.I. H for Henry)

Magnetic field lines:
Stronger field $\leftrightarrow$ higher density of field lines.
 $B_1 = \mu n_1 I_1$

$$n_1 \equiv \frac{N_1}{l_1}$$

inductance.
 $\rightarrow$ Use correct sign for the mutual inductance terms.

Mutual inductance in our equation:

If Solenoids 1 & 2 are ∞ far apart:

$$V_1 = \frac{d\Phi_1}{dt} = \underbrace{(\mu n_1 S_1)}_{L_1 \rightarrow \text{self-inductance}} \frac{dI_1}{dt} \Leftrightarrow \boxed{V_1 = L_1 \frac{dI_1}{dt}}$$

Faraday's Law or Law of induction: $\mathcal{E} = \frac{d\Phi}{dt}$
(a change of magnetic flux Φ in time induces an e.m.f. or voltage)

$\Phi = \oint \vec{B} \cdot d\vec{S}$ For our solenoid: $\vec{B}$ is uniform throughout cross-sectional area & perpendicular to these area. $\rightarrow \Phi_1 = B_1 \cdot S_1$

$$= \mu n_1 I_1 S_1$$

$$\left\{ \begin{array}{l} V_1 = L_1 \frac{dI_1}{dt} \\ V_2 = L_2 \frac{dI_2}{dt} \end{array} \right.$$

If the solenoids are not ∞ far apart: $\left\{ \begin{array}{l} \text{change of current} \\ \text{in \#1 will affect} \\ \text{voltage in \#2} \\ \text{\& vice versa.} \end{array} \right.$

$$\left\{ \begin{array}{l} V_1 = L_1 \frac{dI_1}{dt} + M \frac{dI_2}{dt} \\ V_2 = L_2 \frac{dI_2}{dt} + M \frac{dI_1}{dt} \end{array} \right.$$

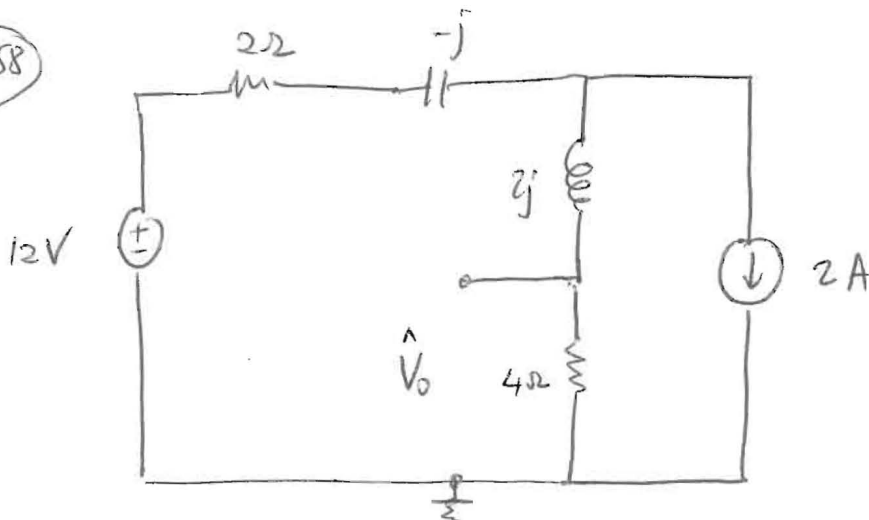
Magnetic coupling (more than one solenoid in a circuit)

$\rightarrow$ Use modified equation for voltages to include mutual inductance.

$\rightarrow$ Use correct sign for the mutual inductance terms.

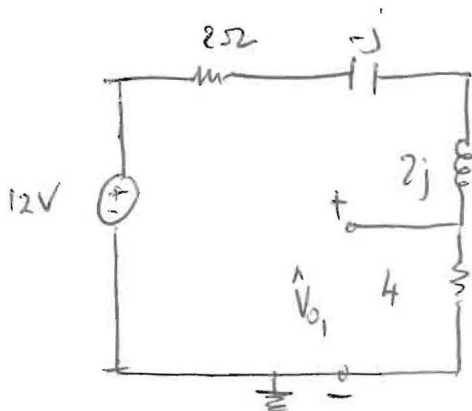
$\rightarrow$ Use the DOT Convention.

7-58

Find $\hat{V}_o$ using superposition.

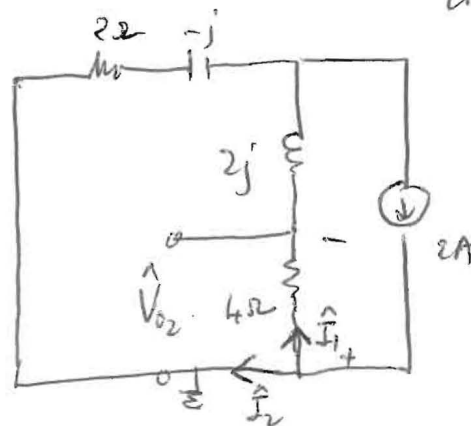
$$\hat{V}_o = \hat{V}_{o1} + \hat{V}_{o2}$$

$\downarrow$ $\downarrow$
 Current source is open-circuited Voltage source is short-circuited.



Voltage Division:

$$\hat{V}_{o1} = 12 \frac{4}{4 + 2 + j} = \frac{48}{6 + j}$$



Current division:

$$\hat{V}_{o2} = -4\hat{I}_1$$

$$\hat{I}_1 = 2A \frac{2 - j}{2 - j + 4 + 2j} = \frac{4 - 2j}{6 + j}$$

$$\hat{V}_o = \frac{48}{6 + j} - \frac{16 - 8j}{6 + j} = \frac{32 + 8j}{6 + j} = 8 \frac{4 + j}{6 + j}$$

$$= 8 \frac{\sqrt{17} \angle \tan^{-1} \frac{1}{4}}{\sqrt{37} \angle \tan^{-1} \frac{1}{6}} = 5.42 \angle 4.53^\circ \checkmark$$

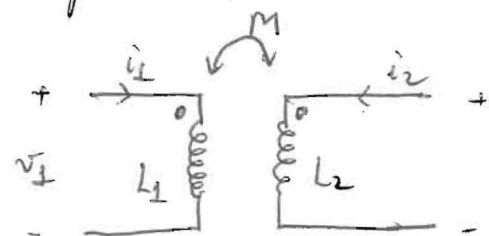
$$V_o(t) = 5.42 \cos(377t + 4.53^\circ) \text{ V}$$

Self & Mutual Inductance: Dot Convention:

Provides information about internal properties of the inductors:
how the wire is wrapped around solenoid determines direction of the magnetic field. $\leftrightarrow$ a reference for correct signs for the mutual inductance terms.

3 possible scenarios involving 2 inductors:

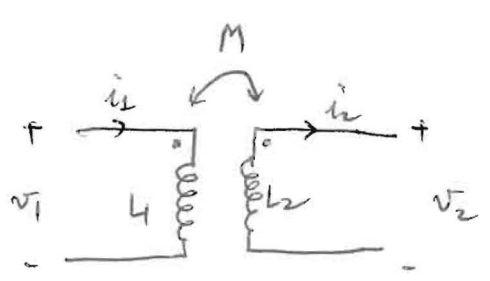
1)



$$\Leftrightarrow \begin{cases} v_1 = L_1 \frac{di_1}{dt} + M \frac{di_2}{dt} \\ v_2 = L_2 \frac{di_2}{dt} + M \frac{di_1}{dt} \end{cases} \quad (\text{Time-domain})$$

$$\begin{cases} \hat{V}_1 = j\omega L_1 \hat{I}_1 + j\omega M \hat{I}_2 \\ \hat{V}_2 = j\omega L_2 \hat{I}_2 + j\omega M \hat{I}_1 \end{cases} \quad (\text{Frequency-domain})$$

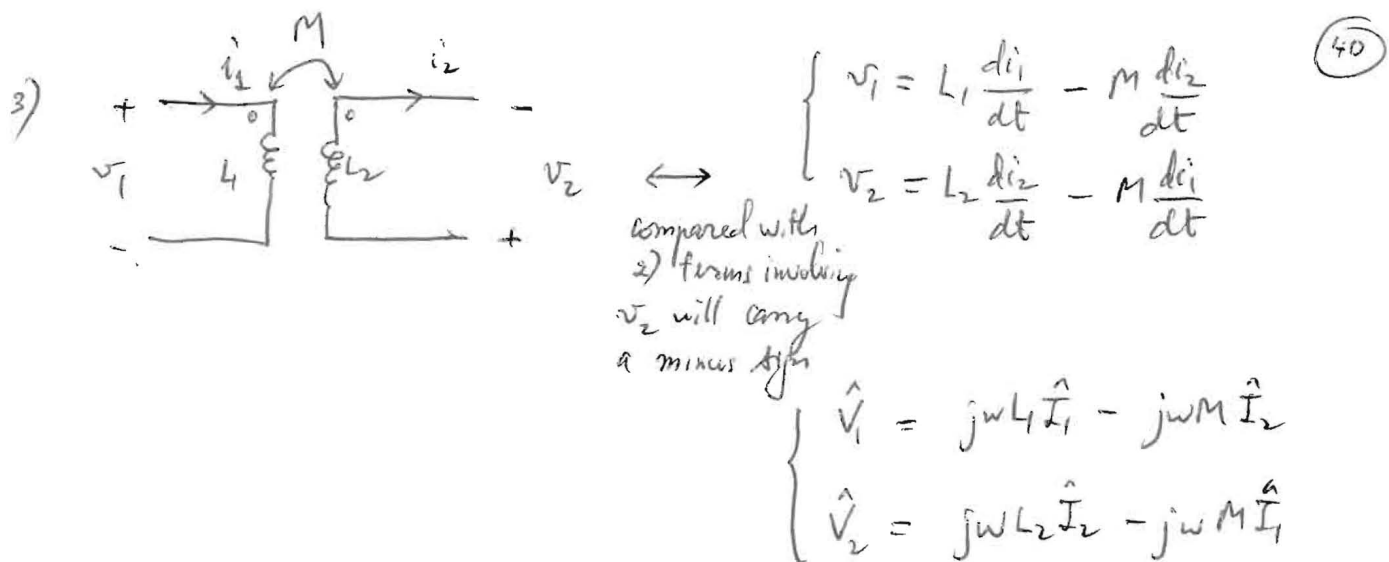
2)



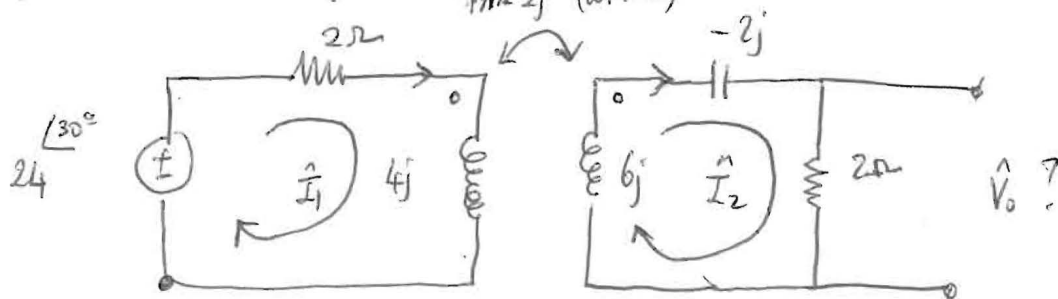
$$\Leftrightarrow \begin{cases} v_1 = L_1 \frac{di_1}{dt} - M \frac{di_2}{dt} \\ v_2 = -L_2 \frac{di_2}{dt} + M \frac{di_1}{dt} \end{cases}$$

terms involving i_2 will carry a minus sign

$$\begin{cases} \hat{V}_1 = j\omega L_1 \hat{I}_1 - j\omega M \hat{I}_2 \\ \hat{V}_2 = -j\omega L_2 \hat{I}_2 + j\omega M \hat{I}_1 \end{cases}$$



Example: AC Analysis of circuits with more than one inductor:



Loop Analysis: choose direction for current in each loop (2 loops)
 $\rightarrow$ CW $\hat{I}_1$ & $\hat{I}_2 \rightarrow$ 2nd scenario! $\rightarrow$ Mutual inductance term has the opposite sign as the self inductance term.

Left loop: 1) $24 \angle 30^\circ - \hat{I}_1 (2 + 4j) + \hat{I}_2 \underbrace{2j}_M = 0$

Right loop: 2) $-\hat{I}_2 (\underbrace{6j}_{L_2} - 2j + 2) + 2j \hat{I}_1 = 0$

$\hat{I}_1 = \hat{I}_2 \frac{(4j + 2)}{2j} = \hat{I}_2 \frac{2j + 1}{j} = \hat{I}_2 (+2 - j)$

1) $24 \angle 30^\circ - \hat{I}_2 (2 - j)(2 + 4j) + \hat{I}_2 2j = 0$
 $24 \angle 30^\circ - \hat{I}_2 [(2 - j)(2 + 4j) - 2j] = 0$

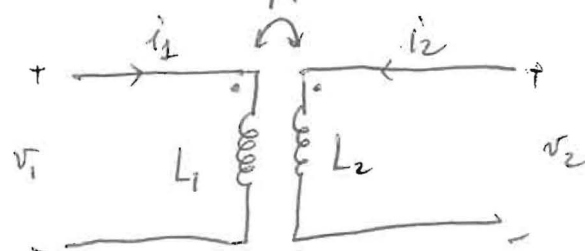
$$24 \angle 30^\circ - \hat{I}_2 [8 + 6j - 2j] = 0$$

$$\hat{I}_2 = \frac{24 \angle 30^\circ}{8 + 4j} = \frac{6 \angle 30^\circ}{2 + j}$$

$$= \frac{6 \angle 30^\circ - \tan^{-1} \frac{1}{2}}{\sqrt{5}} = 2.68 \angle 3.43^\circ \text{ A}$$

$$\hat{V}_0 = 2 \hat{I}_2 = 2 \times 2.68 \angle 3.43^\circ \text{ V} = 5.36 \angle 3.43^\circ \text{ V}$$

Energy Analysis of Magnetically Coupled Circuits :

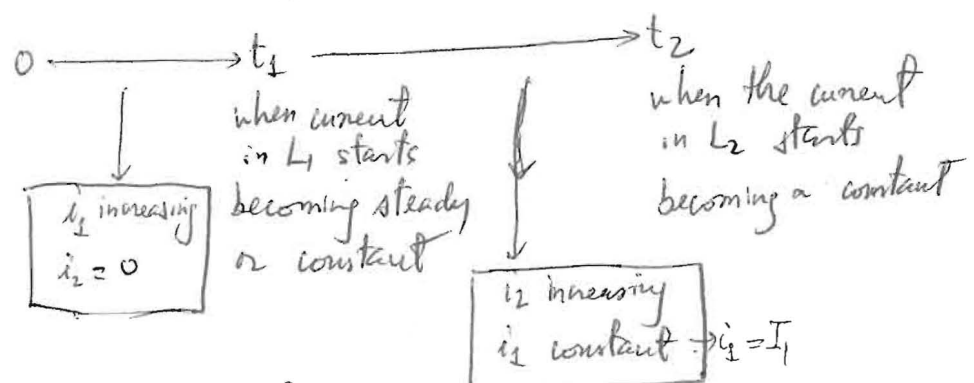


Scenario #1:

$$\begin{cases} v_1(t) = L_1 \frac{di_1}{dt} + M \frac{di_2}{dt} \\ v_2(t) = L_2 \frac{di_2}{dt} + M \frac{di_1}{dt} \end{cases}$$

Power consumption : $P(t) = i_1(t)v_1(t) + i_2(t)v_2(t)$

Energy spent : $\mathcal{E} = \int_0^{t_2} dt P(t)$



$$\mathcal{E} = \int_0^{t_1} dt \underbrace{\left(i_1 L_1 \frac{di_1}{dt} + i_1 M \frac{di_2}{dt} \right)}_{i_1 v_1 \text{ (} i_2 = 0 \text{)}} + \int_{t_1}^{t_2} dt \left[\underbrace{\left(i_1 L_1 \frac{di_1}{dt} + i_1 M \frac{di_2}{dt} \right)}_{i_1 v_1} + \underbrace{\left(i_2 L_2 \frac{di_2}{dt} + i_2 M \frac{di_1}{dt} \right)}_{i_2 v_2} \right]$$

$$\mathcal{E} = \underbrace{L_1 \int_0^{t_1} i_1 di_1}_{\frac{1}{2} L_1 I_1^2} + M I_1 \int_{t_1}^{t_2} di_2 + L_2 \int_{t_1}^{t_2} i_2 di_2$$

$$\mathcal{E} = \boxed{\frac{1}{2} L_1 I_1^2} + \boxed{M I_1 I_2} + \boxed{\frac{1}{2} L_2 I_2^2}$$

$(i_1(t=t_1) = I_1)$ $\downarrow$ Magnetic energy in L_1
 $(i_2(t=t_1) = 0, i_2(t=t_2) = I_2)$ $\downarrow$ Coupling energy b/w L_1 & L_2
 $\downarrow$ Magnetic energy in L_2

$$KE = \frac{1}{2} m v^2$$

$\downarrow$
inertia of motion: mass opposes any ^{instantaneous} change in velocity.

$$ME = \frac{1}{2} L I^2$$

$\downarrow$
inertia of current: inductance opposes any ^{instantaneous} change in current. $\downarrow$ using magnetic energy it stored when current was on

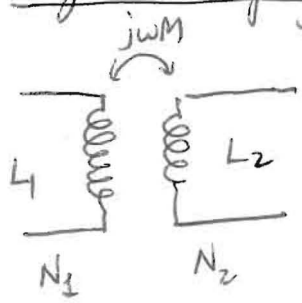
$$EE = \frac{1}{2} C V^2$$

$\downarrow$
inertia of voltage: capacitance opposes any instantaneous change in voltage $\downarrow$ using electric energy it stored when current was on.

inertia $\leftrightarrow$ energy storage

Ideal Transformer:

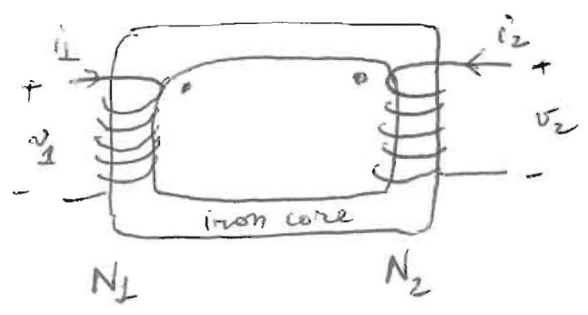
Magnetic coupling:



→ Two solenoids with different inductances L_1 & L_2 (different cross sections, numbers of turns, etc.)

Ideal Transformer:

Two solenoids of same cross sections: using an iron core:



Faraday's Law: voltage v_1 comes from a change of magnetic flux through solenoid #1: $N_1 \Phi_1$

$$v_1 = \frac{d(N_1 \Phi_1)}{dt} = N_1 \frac{d\Phi_1}{dt}$$

$$\Phi_1 = \text{m. flux through one loop of inductor or solenoid \#1}$$

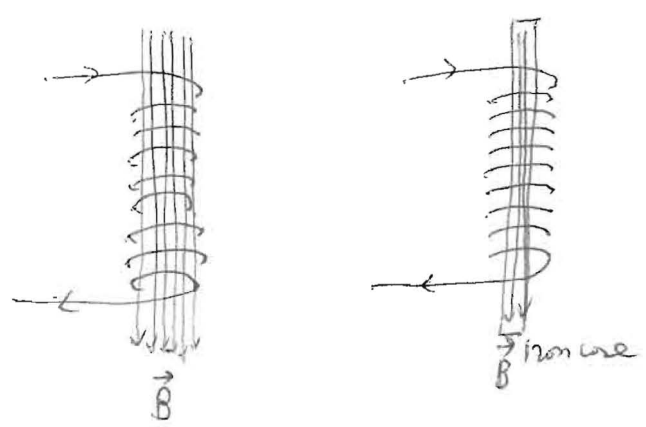
$$= B \cdot S_1$$

$$v_2 = \frac{d(N_2 \Phi_2)}{dt} = N_2 \frac{d\Phi_2}{dt}$$

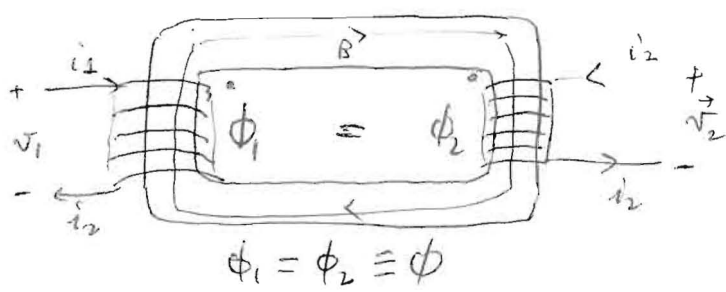
$$\Phi_2 = B \cdot S_2$$

Two inductors or solenoids next to each other

ideal Transformers (same cross section)



can set a fixed cross-section for the magnetic flux by using an iron core



$$v_1 = N_1 \frac{d\Phi_1}{dt} = N_1 \frac{d\Phi}{dt}$$

$$v_2 = N_2 \frac{d\Phi_2}{dt} = N_2 \frac{d\Phi}{dt}$$

$$\frac{v_1}{v_2} = \frac{N_1}{N_2} \text{ voltage transformer.}$$

Any connection b/w i_1 & i_2 in an ideal transformer?

Yes, using Ampere's Law.

$$\oint_{\text{Amp. loop}} \vec{H} \cdot d\vec{l} = \text{current enclosed by loop}$$

$$\vec{B} = \mu \vec{H}$$

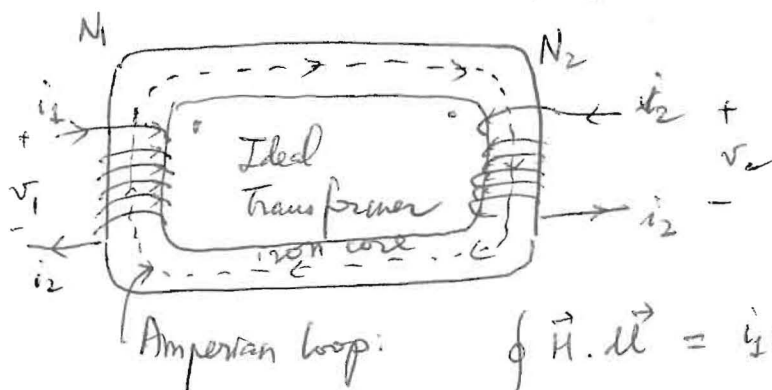
magnetic induction

magnetic field

magnetic permeability

(most of the time we call $\vec{B}$ the magnetic field)

$$\oint_{\text{Amp loop}} \vec{B} \cdot d\vec{l} = \mu I_{\text{enclosed by loop}}$$



Amperian loop:

$$\oint_{\text{A. loop}} \vec{H} \cdot d\vec{l} = i_1 N_1 + i_2 N_2$$

$$\frac{1}{\mu} \oint_{\text{A. loop}} \vec{B} \cdot d\vec{l}$$

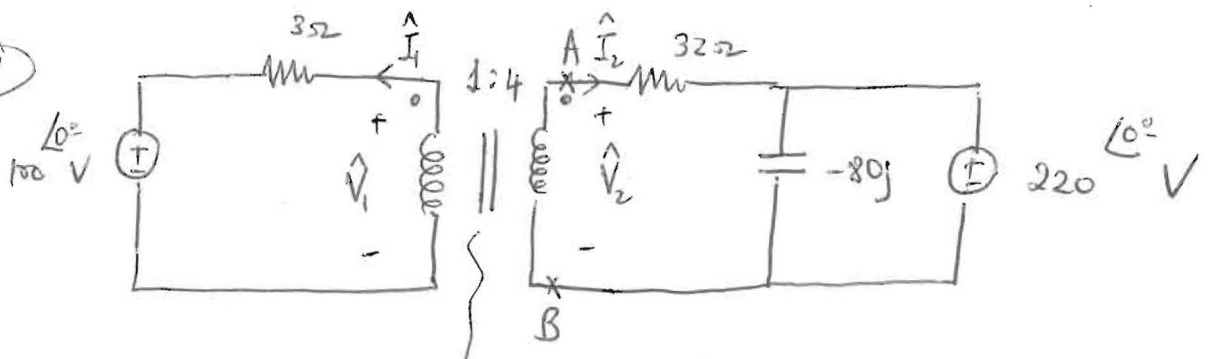
Ideal transformer $\rightarrow$ iron core $= \mu = \infty \rightarrow \frac{1}{\mu_{\text{iron core}}} = 0$

$$\rightarrow 0 = i_1 N_1 + i_2 N_2 \rightarrow \boxed{\frac{i_1}{i_2} = -\frac{N_2}{N_1}}$$

Ideal Transformer.

8.51

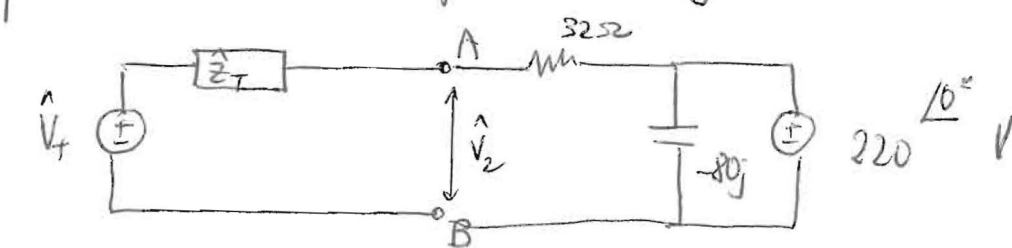
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iron core ideal transformer coil ratio: $\frac{N_1}{N_2} = \frac{1}{4}$
 $\hat{I}_1, \hat{I}_2, \hat{V}_1, \hat{V}_2$?

Ideal Transformer: $\frac{\hat{V}_2}{\hat{V}_1} = 4$; $\frac{\hat{I}_2}{\hat{I}_1} = -\frac{1}{4}$ Only need to solve for two variables, e.g. $\hat{I}_2$ & $\hat{V}_2$

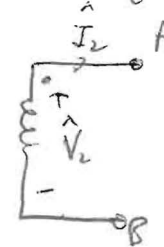
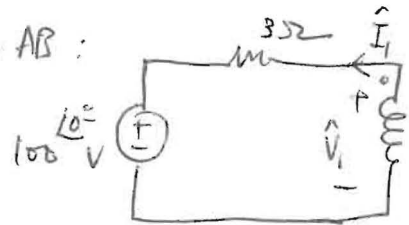
Solve using Thevenin equivalent: $\rightarrow$ b/w two points $\rightarrow$ A & B
 (we need to find $\hat{V}_2$)
 will replace circuit to the left of AB by its Thevenin's equivalent



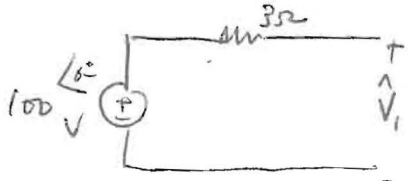
Find this
 $\hat{V}_T = \hat{V}_{oc}$ across AB
 (open circuit)

$\hat{Z}_T =$ impedance b/w A & B when sources are eliminated
 Voltage source $\rightarrow$ short-circuit
 current source $\rightarrow$ open-circuit

$\hat{V}_{oc}$ across AB:

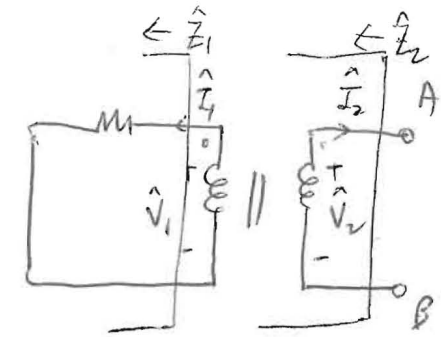


$\hat{V}_{oc} = \hat{V}_2$
 $\hat{I}_2 = 0 \rightarrow \hat{I}_1 = -4\hat{I}_2 = 0$
 $\rightarrow$ Left loop is open!



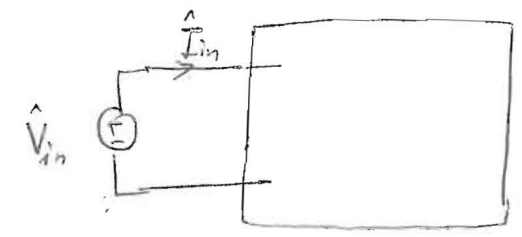
$\rightarrow \hat{V}_1 = 100 \angle 0^\circ \text{ V} \rightarrow \hat{V}_2 = 4\hat{V}_1 = 400 \angle 0^\circ \text{ V}$
 $\rightarrow \boxed{\hat{V}_{oc} = 400 \angle 0^\circ \text{ V} = \hat{V}_T}$

$\hat{Z}_T = \hat{Z}_{AB} / \text{no source}$



$\hat{Z}_{AB} = ? = \frac{\hat{V}_2}{-\hat{I}_2}$

In general the input impedance is found:

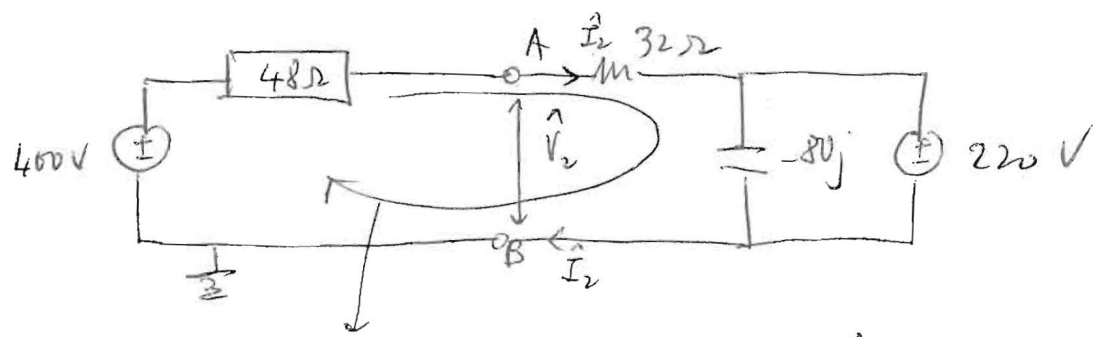


$\hat{Z}_{in} = \frac{\hat{V}_{in}}{\hat{I}_{in}}$

$\hat{Z}_T = \frac{\hat{V}_2}{-\hat{I}_2} = \frac{4\hat{V}_1}{-(-\frac{1}{4}\hat{I}_1)} = 16 \frac{\hat{V}_1}{\hat{I}_1} = 16.3 = 48\Omega$

input impedance from the voltage source of 100 inductor #1 to the left.

16 fold increase due to the ideal transformer.



KVL: $400 - \underbrace{(48 + 32)}_{80} \hat{I}_2 - 220 = 0$

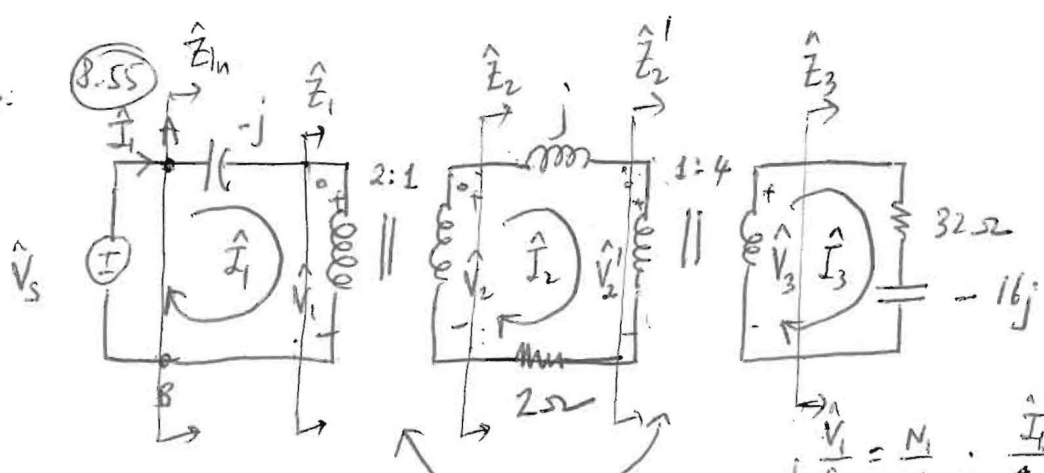
voltage drop across $(-80j)$

$\hat{I}_2 = \frac{400 - 220}{80} = \frac{180}{80} = 2.25A$

$\hat{V}_2 = \hat{V}_{AB} = 400 - 48\hat{I}_2 = 292V \rightarrow \hat{V}_1 = \frac{\hat{V}_2}{4} = \frac{292}{4} = 73V$

$\hat{I}_1 = -4\hat{I}_2 = -4 \times 2.25 = -9A$

HW4:



ideal transformers $\left\{ \begin{array}{l} \frac{\hat{V}_1}{\hat{V}_2} = \frac{N_1}{N_2} \cdot \frac{\hat{I}_1}{\hat{I}_2} = -\frac{N_2}{N_1} \text{ w/ dots at tops} \\ \text{(both currents go down @ inductors)} \end{array} \right.$

Find input impedance b/w A & B from the source $\hat{V}_s$:

$$\hat{Z}_{in} = \frac{\hat{V}_s}{\hat{I}_1} \quad (\text{low } \hat{I}_1 \Rightarrow \text{high } \hat{Z}_{in} \text{ as in meters})$$

Strategy: find $\hat{I}_1$ in term of $\hat{V}_s$

- 1) Name voltages across inductors as part of the ideal transformers.
Name currents in each of the three loops: $\hat{V}_1, \hat{V}_2, \hat{V}_2', \hat{V}_3$
 $\hat{I}_1, \hat{I}_2, \hat{I}_3$

- 2) Starting w/ loop #3:

$$\frac{\hat{V}_3}{\hat{I}_3} = \hat{Z}_3 = 32 - 16j \quad \rightarrow \quad \frac{4\hat{V}_2'}{\hat{I}_2} = 32 - 16j$$

$$\left. \begin{array}{l} \text{Since } \frac{\hat{V}_3}{\hat{V}_2'} = 4 \rightarrow \hat{V}_3 = 4\hat{V}_2' \\ \frac{-\hat{I}_3}{\hat{I}_2} = -\frac{1}{4} \rightarrow \hat{I}_3 = \frac{\hat{I}_2}{4} \end{array} \right\}$$

$$\Rightarrow \frac{\hat{V}_2'}{\hat{I}_2} = \hat{Z}_2' = 2 - j$$

- 3) $\hat{Z}_2 = 4 \checkmark$ (series combination of j & $2-j$ & 2)
 $\hat{Z}_2 = \frac{\hat{V}_2}{\hat{I}_2}$ how does $\hat{V}_2$ relate to $\hat{V}_2'$?

KVL for middle loop: $-\hat{V}_2 + \hat{I}_2 j + \hat{V}_2' + 2\hat{I}_2 = 0$

$$\hat{V}_2 = (2+j)\hat{I}_2 + \hat{V}_2' = \hat{I}_2 \left[2+j + \frac{\hat{V}_2'}{\hat{I}_2} \right] = 4\hat{I}_2$$

$$\rightarrow \hat{Z}_2 = \frac{\hat{V}_2}{\hat{I}_2} = 4 \checkmark$$

4) Find $\hat{Z}_1 = \frac{\hat{V}_1}{\hat{I}_1} = \frac{2\hat{V}_2}{\frac{\hat{I}_2}{2}} = 4 \frac{\hat{V}_2}{\hat{I}_2} = 4\hat{Z}_2 = 4 \times 4 = 16 \Omega$

Ideal Transformer $\frac{\hat{V}_1}{\hat{V}_2} = 2 \rightarrow \hat{V}_1 = 2\hat{V}_2$
 $\frac{-\hat{I}_2}{\hat{I}_1} = -2 \rightarrow \hat{I}_1 = \frac{\hat{I}_2}{2}$

$\hat{Z}_1 = \frac{\hat{V}_1}{\hat{I}_1} = 16$

5) Find $\hat{Z}_{in} =$

Method #1: series combination of $-j$ & $\hat{Z}_1 = -j + \hat{Z}_1 = -j + 16 \Omega$

Method #2: $\hat{Z}_{in} = \frac{\hat{V}_s}{\hat{I}_1} = \frac{\hat{V}_1 - \hat{I}_1 j}{\hat{I}_1} = \frac{\hat{V}_1}{\hat{I}_1} - j = 16 - j \Omega$

Relate $\hat{V}_s$ to $\hat{V}_1$: $\hat{V}_s = \hat{V}_1 + (-j\hat{I}_1)$
(KVL on loop #1)

Power Analysis:

- What load would extract the most power?
- What is the average power consumed in an AC circuit?
- What are the root-mean-square (RMS) or effective values for current & voltages?
- What is the complex power?

Instantaneous power: p (lower case letter p)

$$p(t) = i(t) \cdot v(t) = I_m \cos(\omega t + \theta_i) \cdot V_m \cos(\omega t + \theta_v)$$

AC currents:
$$\left. \begin{aligned} i(t) &= I_m \cos(\omega t + \theta_i) \\ v(t) &= V_m \cos(\omega t + \theta_v) \end{aligned} \right\} \begin{array}{l} I_m \text{ \& } V_m \text{ magnitude of} \\ \text{alternating current \& } \\ \text{voltage.} \\ \theta_i \text{ \& } \theta_v \text{ are phases of} \\ \text{AC current \& voltage.} \end{array}$$

i & v are time varying time independent.

Average power: P (upper case letter P)

$$P = \frac{1}{T} \int_{t_0}^{t_0+T} dt \cdot p(t) \quad \left(\text{integrating the instantaneous power over one period, then divide by the period } T \right)$$

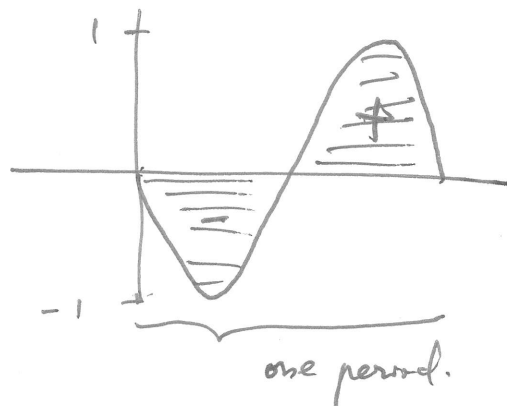
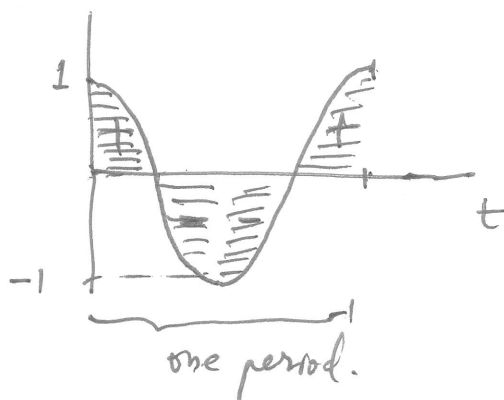
Any connection b/w T & ω : $T = \frac{2\pi}{\omega}$

Need to integrate a product of 2 cosines: $\rightarrow$

$$\left\{ \begin{aligned} \cos \alpha \cos \beta &= \frac{1}{2} [\cos(\alpha - \beta) + \cos(\alpha + \beta)] \quad \text{Trig. Id} \\ \alpha &\equiv \omega t + \theta_i \rightarrow \alpha - \beta = \theta_i - \theta_v; \alpha + \beta = 2\omega t + \theta_i + \theta_v \\ \beta &\equiv \omega t + \theta_v \end{aligned} \right.$$

$$\begin{aligned}
 \rightarrow \underline{P} &= \frac{1}{T} I_m V_m \frac{1}{2} \left[\int_T dt \underbrace{\cos(\theta_i - \theta_v)}_{\text{time independent}} + \int_T dt \underbrace{\cos(2\omega t + \theta_i + \theta_v)}_{\text{period is } \frac{2\pi}{2\omega} = \frac{T}{2}} \right] \\
 &= \frac{1}{2} I_m V_m \frac{1}{T} \left[\underbrace{\cos(\theta_i - \theta_v)}_{\text{time independent}} \underbrace{\int_T dt}_{T} + \underbrace{\int_T dt \cos(2\omega t + \theta_i + \theta_v)}_{\text{integral of a sinusoid over two of its periods of } \frac{T}{2}} \right] \\
 &\quad \Downarrow 0
 \end{aligned}$$

* Integral of a sinusoid over one of its period is: 0



$$\Rightarrow \underline{P} = \frac{1}{2} I_m V_m \cos(\theta_i - \theta_v)$$

Average Power is not time-varying!

Consequences:

1) At a pure resistive element: $\theta_i = \theta_v$
(a resistor)
 $\rightarrow$ max. ave. power: $\underline{P}_{\max} = \frac{1}{2} I_m V_m$

2) At a pure reactive element: $\theta_i - \theta_v = \pm \frac{\pi}{2}$
(inductor or a capacitor)

$$\underline{P} = 0$$

(they are energy storage devices $\rightarrow$ makes sense)

At a resistor

$$\hat{Z} = \frac{\hat{V}}{\hat{I}} = \frac{V_m \angle \theta_v}{I_m \angle \theta_i} = \frac{V_m}{I_m} \angle \theta_v - \theta_i$$

Ohm's law works
with phasors

because we know $\hat{Z} = R \rightarrow$ a real number $\rightarrow$ phase is
zero $\rightarrow \theta_v - \theta_i = 0$ or $\boxed{\theta_v = \theta_i}$

At an inductor

$$\hat{Z}_L = \frac{\hat{V}}{\hat{I}} = j\omega L = \omega L \angle 90^\circ$$

$$= \frac{V_m \angle \theta_v}{I_m \angle \theta_i} \quad \left. \vphantom{\frac{V_m \angle \theta_v}{I_m \angle \theta_i}} \right\} \boxed{\theta_v - \theta_i = 90^\circ}$$

At a capacitor:

$$\hat{Z}_C = \frac{1}{j\omega C} = \frac{1}{\omega C} \angle -90^\circ$$

$$= \frac{V_m \angle \theta_v}{I_m \angle \theta_i} \quad \left. \vphantom{\frac{V_m \angle \theta_v}{I_m \angle \theta_i}} \right\} \boxed{\theta_v - \theta_i = -90^\circ}$$

What are the R-M-S values or effective values for current & voltage?
(Root of the Mean of the square)

$$I_{RMS} = \sqrt{\frac{1}{T} \int_T dt \, i^2(t)}$$

mean of square of
current.

AC current: $i(t) = I_m \cos(\omega t + \theta_i)$

$$I_{RMS} = I_m \sqrt{\frac{1}{T} \int_T dt \cos^2(\omega t + \theta_i)}$$

Trig. Id: $\cos^2 \alpha = \frac{1}{2}(1 + \cos 2\alpha)$

$$I_{RMS} = \frac{I_m}{\sqrt{2}} \sqrt{\frac{1}{T} \int_T dt (1 + \cos(2\omega t + 2\theta_i))} = \frac{I_m}{\sqrt{2}} \sqrt{\underbrace{\frac{1}{T} \int_T dt}_1 + \underbrace{\frac{1}{T} \int_T dt \cos(2\omega t + 2\theta_i)}_0}$$

period of $\frac{T}{2}$

$$\boxed{I_{RMS} = \frac{I_m}{\sqrt{2}}}$$

likewise the RMS value for voltage is its magnitude divided by $\sqrt{2}$:

$$V_{RMS} = \frac{V_m}{\sqrt{2}}$$

Consequences:

1) Average Power in term of the RMS values:

$$\begin{aligned} P &= \frac{1}{2} I_m V_m \cos(\theta_i - \theta_v) \\ &= \frac{I_m}{\sqrt{2}} \cdot \frac{V_m}{\sqrt{2}} \cdot pf \end{aligned} \quad \left\{ \begin{array}{l} P = I_{RMS} \cdot V_{RMS} \cdot pf \end{array} \right.$$

Power factor: pf (always < 1)
 $pf = \cos(\theta_i - \theta_v)$

2) Pure-resistive element: $pf = \cos(0) = 1 \rightarrow P = I_{RMS} \cdot V_{RMS}$
 $\theta_i = \theta_v$ (reminds us about DC analysis)

Pure-reactive element: $pf = \cos(\pm \frac{\pi}{2}) = 0 \rightarrow P = 0$
 $\theta_i - \theta_v = \pm \frac{\pi}{2}$

Complex Power: upper case letter $\hat{S}$

Not a measurable quantity, but a definition that facilitates the analysis of power in AC circuits (similar to the complex phasors for the analysis of AC circuits)

$$\hat{S} \equiv \hat{V}_{RMS} \cdot \hat{I}_{RMS}^* = V_{RMS} \angle \theta_v \cdot I_{RMS} \angle -\theta_i = V_{RMS} \cdot I_{RMS} \angle \theta_v - \theta_i$$

Complex conjugate

$$= \underbrace{V_{RMS} \cdot I_{RMS}}_{\substack{\uparrow \\ \text{de Moivre's formula}}} \cos(\theta_v - \theta_i) + j V_{RMS} \cdot I_{RMS} \sin(\theta_v - \theta_i)$$

de Moivre's formula

The real part of the complex power is just

$$I_{RMS} \cdot V_{RMS} \text{ pf}$$

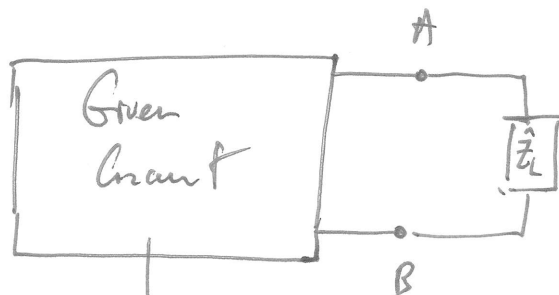
$$(\cos(\theta_v - \theta_i) = \cos(\theta_i - \theta_v) = \text{pf})$$

since $\cos$ is symmetric w/r/t the origin or $\cos$ is an even function)

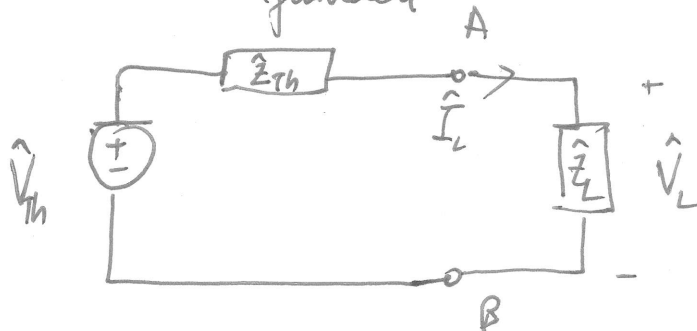
which is the average power P .

$$\Rightarrow \hat{S} \equiv \underbrace{P}_{\text{Real part}} + j \underbrace{V_{RMS} I_{RMS} \sin(\theta_v - \theta_i)}_{\text{Imaginary part}}$$

Which load would draw the max power given a circuit?



Theremin's equivalent



what $\hat{Z}_L$ will draw the most power out of this given circuit?

Average power consumed at load is P_L

Max P_L can be achieved when

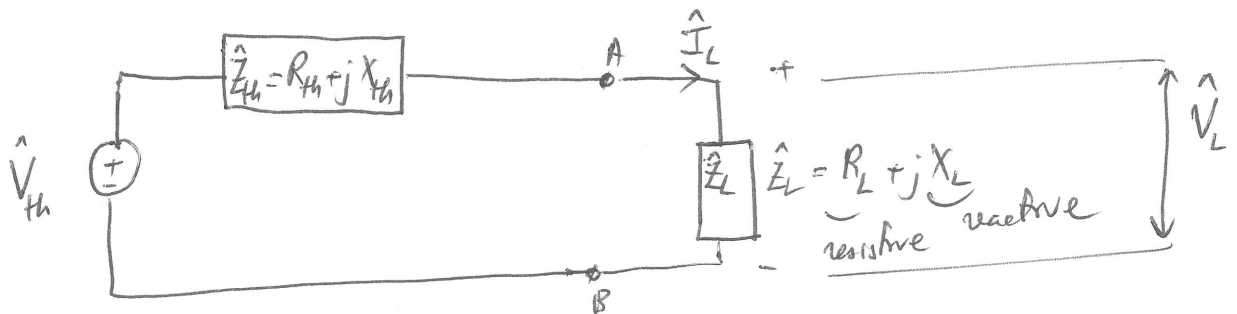
$$\hat{Z}_L = \hat{Z}_{th}^*$$

Proof of the Max Power Theorem:

In general, for the circuit shown in the previous figure, the average power consumed at the load is $P_L = \frac{1}{2} I_L V_L \cos(\theta_V - \theta_I)$ (time-independent).

Using the Thevenin's equivalent for the given circuit, we can write I_L & V_L in terms of V_{th} , R_{th} , X_{th} where

$$V_{th} = |\hat{V}_{th}|; \quad \hat{Z}_{th} = \underbrace{R_{th}}_{\text{resistive}} + j \underbrace{X_{th}}_{\text{reactive}}$$



$$a) \quad \hat{I}_L = \frac{\hat{V}_{th}}{\hat{Z}_{th} + \hat{Z}_L} \rightarrow I_L \quad (\text{to use in } P_L) = \frac{V_{th}}{\sqrt{(R_{th} + R_L)^2 + (X_{th} + X_L)^2}}$$

(real part & imaginary part are perpendicular in the complex plane) $\rightarrow$ can use Pythagoras Theorem to find the magnitude of the total impedance.

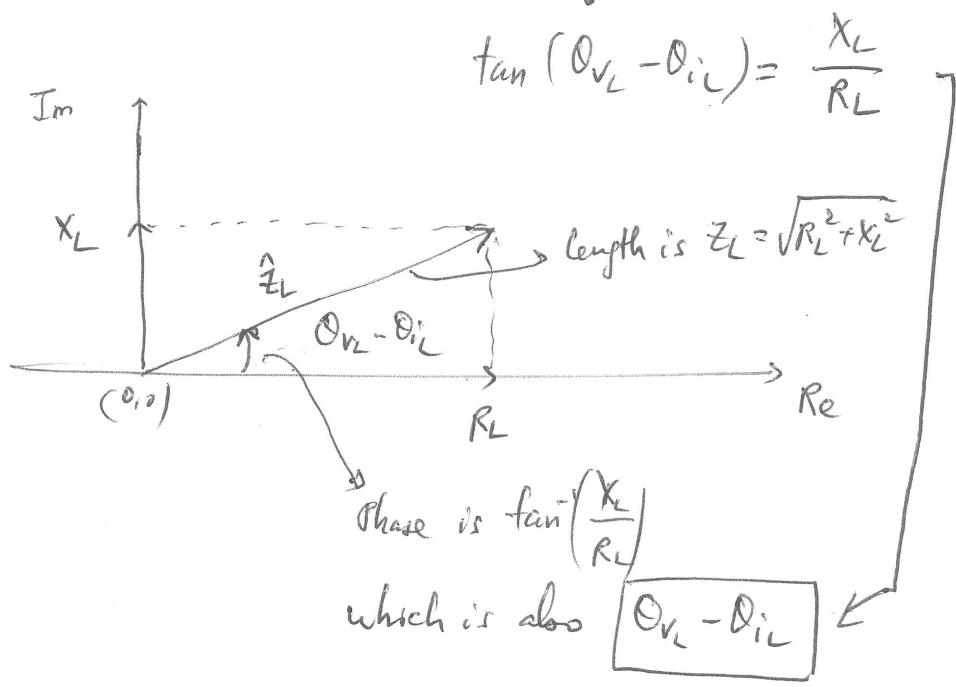
$$\hat{V}_L = \hat{V}_{th} \frac{\hat{Z}_L}{\hat{Z}_L + \hat{Z}_{th}}$$

$$\rightarrow V_L \quad (\text{to be used in } P_L) = \frac{V_{th} \sqrt{R_L^2 + X_L^2}}{\sqrt{(R_{th} + R_L)^2 + (X_{th} + X_L)^2}}$$

$$b) \quad \cos(\theta_{V_L} - \theta_{I_L})$$

$$\left[\hat{Z}_L = \frac{\hat{V}_L}{\hat{I}_L} = \frac{V_L}{I_L} \angle \theta_{V_L} - \theta_{I_L} \right] \rightarrow \hat{Z}_L = R_L + jX_L = \sqrt{R_L^2 + X_L^2} \angle \tan^{-1} \frac{X_L}{R_L}$$

$\downarrow$ Ohm's Law $\downarrow$ Cartesian format $\downarrow$ Polar format



Then: $\cos(\theta_{V_L} - \theta_{I_L}) = \frac{R_L}{\sqrt{R_L^2 + X_L^2}}$

Then plug a) & b) into $P_L = \frac{1}{2} I_L V_L \cos(\theta_{V_L} - \theta_{I_L})$

$$P_L = \frac{1}{2} \underbrace{\frac{V_{Th}}{\sqrt{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2}}}_{I_L} \cdot \underbrace{\frac{V_{Th} \sqrt{R_L^2 + X_L^2}}{\sqrt{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2}}}_{V_L} \cdot \underbrace{\frac{R_L}{\sqrt{R_L^2 + X_L^2}}}_{\cos(\theta_{V_L} - \theta_{I_L})}$$

$$P_L = \frac{1}{2} \frac{V_{Th}^2 R_L}{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2}$$

Ave. power @ load in terms of the Thevenin equivalent for circuit & load information.

What X_L & R_L will maximize P_L ? Use Calculus to find the maximum of a function:

- i) Larger P_L needs smaller denominator: if we drop a parenthesis (since they are both squared or +)

$$\rightarrow (X_{th} + X_L)^2 = 0 \quad \text{or} \quad X_{th} + X_L = 0 \quad \text{or} \quad \boxed{X_L = -X_{th}}$$

This can be realized in practice: we should be able to select a load whose reactive part of the impedance is the same & opposite of that of the Thevenin impedance of the given circuit.

- ii) Then what R_L will maximize P_L ? $\rightarrow$ Use Calculus:

$$\frac{\partial P_L}{\partial R_L} = 0$$

$$\frac{\partial \left[\frac{1}{2} V_{th}^2 \frac{R_L}{(R_{th} + R_L)^2} \right]}{\partial R_L} = \frac{1}{2} V_{th}^2 \frac{\partial \left[\frac{R_L}{(R_{th} + R_L)^2} \right]}{\partial R_L} = \frac{1}{2} V_{th}^2 \left\{ \frac{1}{(R_{th} + R_L)^2} - \frac{2R_L}{(R_{th} + R_L)^3} \right\}$$

has to be 0.

$$\frac{1}{(R_L + R_{th})^2} \left\{ 1 - \frac{2R_L}{R_{th} + R_L} \right\}$$

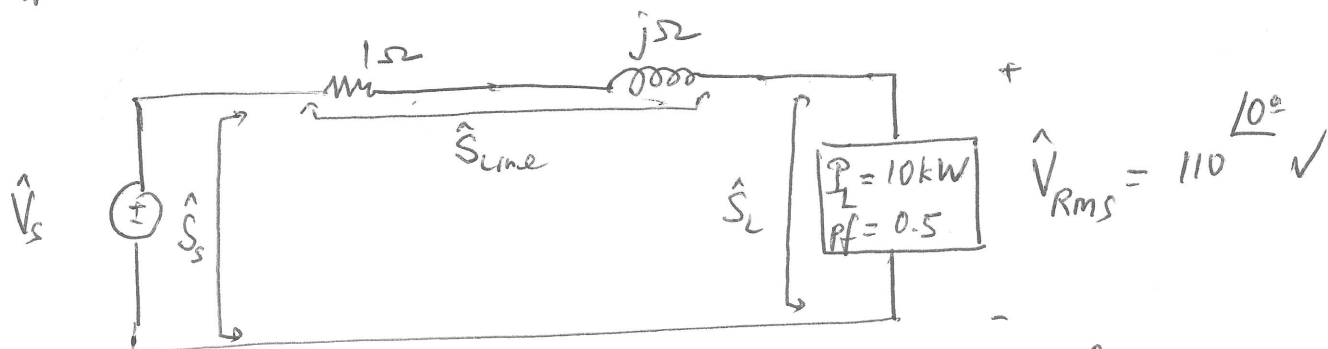
has to be 0 $\rightarrow 2R_L = R_{th} + R_L$

$$\boxed{R_L = R_{th}}$$

Summary:

$$\begin{matrix} R_L = R_{th} \\ X_L = -X_{th} \end{matrix} \rightarrow \boxed{\hat{Z}_L = \hat{Z}_{th}^*}$$

Application of Complex Power $\hat{S}$: scenarios in which $\hat{S}$ is helpful



Calculate $\hat{V}_s$ so the circuit can deliver the specified load.

Strategy using complex power: 1) & 2)

1) Find $\hat{S}_s$ (complex power delivered at source):
 $\hat{S}_s \equiv \hat{V}_s \hat{I}_s^*$ since can then find 2) $\hat{V}_s = \frac{\hat{S}_s}{\hat{I}_s^*}$ (and $\hat{I}_s = \hat{I}_L$)

2) $\hat{S}_s = \hat{S}_{line} + \hat{S}_L$ (power delivered is to be consumed along the line & at the load)

Find $\hat{S}_L$: Recall: $\text{Re}[\hat{S}_L] = P_L = 10\text{kW}$
 Also: $\hat{S}_L = \hat{V}_L \hat{I}_L^* = \frac{V_L}{\text{RMS}} \frac{I_L}{\text{RMS}} \angle \theta_{V_L} - \theta_{I_L}$
 or the angle of $\hat{S}_L$ is simply $\theta_{V_L} - \theta_{I_L}$ (phase)
 $\equiv \cos^{-1}(\text{pf})$

($\text{pf} \equiv \cos(\theta_V - \theta_I)$)

Now what is S_L ? (magnitude of the complex power $\hat{S}_L$)

It is $V_{\text{RMS}_L} I_{\text{RMS}_L} = \frac{P_L}{\text{pf}}$ (since $P = V_{\text{RMS}} I_{\text{RMS}} \text{pf}$)

$$\hat{S}_L = S_L \angle \theta_{V_L} - \theta_{I_L} = \frac{P_L}{\text{pf}} \angle \cos^{-1} \text{pf} = \frac{10000}{0.5} \angle \cos^{-1} 0.5$$

$$\hat{S}_L = 20000 \angle 60^\circ \text{ V.A} \rightarrow \text{save W for real power (average power)}$$

Complex power @ load can be obtain from its average power & pf

$$\hat{S}_L = \frac{P_L}{\text{pf}} \angle \cos^{-1} \text{pf}$$

1 ii) Find $\hat{S}_{\text{Line}}$

$$\hat{S}_{\text{Line}} = \hat{V}_{\text{Line RMS}} \hat{I}_{\text{Line RMS}}^* = \hat{Z}_{\text{Line}} \underbrace{\hat{I}_{\text{Line RMS}} \cdot \hat{I}_{\text{Line RMS}}^*}_{I_{\text{Line RMS}}^2} = I_{\text{Line RMS}}^2 \hat{Z}_{\text{Line}} = I_{\text{Line RMS}}^2 (1+j)$$

Ohm's law
for phasors

$$\hat{S}_{\text{Line}} = I_{\text{Line RMS}}^2 (1+j)$$

Now $I_{\text{Line RMS}} = I_L$ for this circuit.

Recall: $P_L = I_{L \text{ RMS}} \underbrace{V_{L \text{ RMS}}}_{\substack{\downarrow \\ \text{given as } 110}} \text{ pf} \rightarrow I_{L \text{ RMS}} = \frac{P_L}{V_{L \text{ RMS}} \text{ pf}} = \frac{10000}{110 \times 0.5}$

$$I_{L \text{ RMS}} = \frac{20000}{110} = 181.82 \text{ A}$$

$$\hat{S}_{\text{Line}} = 181.82^2 (1+j)$$

1 iii) $\hat{S}_S = \hat{S}_{\text{Line}} + \hat{S}_L = 181.82^2 (1+j) + 20000 (\cos 60^\circ + j \sin 60^\circ)$

$$= (181.82^2 + 20000 \cos 60^\circ) + j (181.82^2 + 20000 \sin 60^\circ)$$

$$= 43058.5 + j 50379$$

$$\hat{S}_S = 66272.8 \angle 49.5^\circ \text{ V-A}$$

(59)

$$2) \left[\hat{V}_S = \frac{\hat{S}_S}{\hat{I}_S^*} = \frac{66272.8 \angle 49.5^\circ}{181.82 \angle 60^\circ} = 364.49 \angle -10.5^\circ \text{ V} \right]$$

$$\hat{I}_S = \hat{I}_{\text{line}} = \hat{I}_L \text{ for our circuit}$$

$$\hat{I}_L^* = \frac{\hat{S}_L}{\hat{V}_L} = \frac{20000 \angle 60^\circ}{110 \angle 0^\circ} = 181.82 \angle 60^\circ$$

$$pf_{\text{source}} = \cos(\text{phase of } \hat{S}_S) = \cos(49.5^\circ) = 0.65$$
