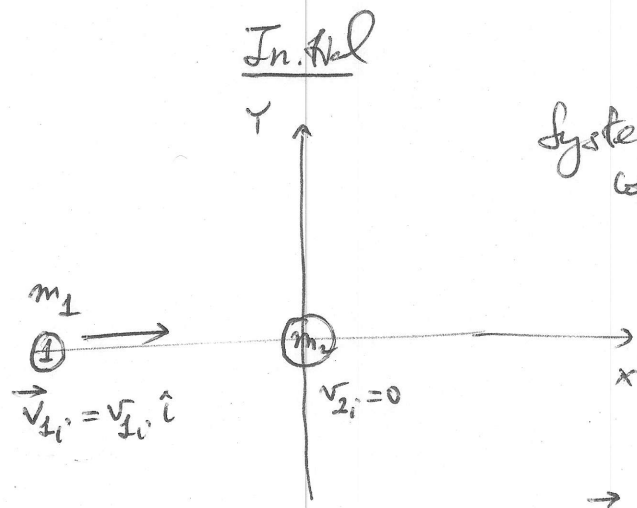
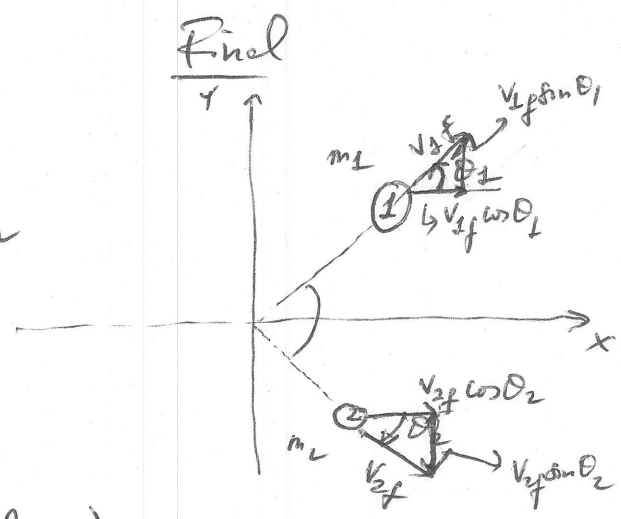


Elastic Collision in 2D:

We will prove that the final directions of two colliding objects (elastic collisions: no deformation, hard-ball collisions) of equal masses form 90° angle!



System of 2 components m_1 & m_2



Elastic collisions:

$$\begin{cases} \vec{P}_i = \vec{P}_f & (2 \text{ eqs}) \\ KE_i = KE_f & (1 \text{ eq.}) \end{cases}$$

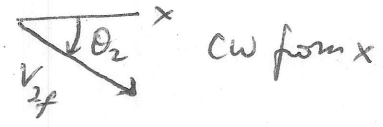
Known variables: $m_1, m_2, \vec{v}_{1i}, \theta_1$ (final direction for m_1)
 ↓
 can only solve for 3 unknowns
 (v_{1f}, v_{2f}, θ_2)

Write the conservation laws in terms of the 2 components of the system.

$$\begin{aligned} \textcircled{1} \quad m_1 v_{1i} &= m_1 v_{1f} \cos \theta_1 + m_2 v_{2f} \cos \theta_2 \\ \textcircled{2} \quad 0 &= m_1 v_{1f} \sin \theta_1 + m_2 v_{2f} \sin \theta_2 \end{aligned}$$

↳ Constant of θ_2 is negative as shown

Standard convention for angle: } + CCW from x-axis } - CW from x-axis.



Eq ② is also: $0 = m_1 v_{1f} \sin \theta_1 - m_2 v_{2f} \sin \theta_2$

③ $\frac{1}{2} m_1 v_{1i}^2 = \frac{1}{2} m_1 v_{1f}^2 + \frac{1}{2} m_2 v_{2f}^2$

Now use math manipulation (algebra) to solve for v_{1f}, v_{2f}, θ_2

①² $m_1^2 v_{1i}^2 = m_1^2 \left(v_{1f}^2 \cos^2 \theta_1 + \frac{m_2^2}{m_1} v_{2f}^2 \cos^2 \theta_2 \right)^2$
 $v_{1i}^2 = \boxed{v_{1f}^2 \cos^2 \theta_1} + \frac{m_2^2}{m_1^2} v_{2f}^2 \cos^2 \theta_2 + 2 \frac{m_2}{m_1} v_{1f} v_{2f} \cos \theta_1 \cos \theta_2$

②² $0 = m_1^2 \left(v_{1f} \sin \theta_1 + \frac{m_2}{m_1} v_{2f} \sin \theta_2 \right)^2$
 $0 = \boxed{v_{1f}^2 \sin^2 \theta_1} + \frac{m_2^2}{m_1^2} v_{2f}^2 \sin^2 \theta_2 + 2 \frac{m_2}{m_1} v_{1f} v_{2f} \sin \theta_1 \sin \theta_2$

$\cos^2 \alpha + \sin^2 \alpha = 1$

①² + ②²: $v_{1i}^2 = v_{1f}^2 + \frac{m_2^2}{m_1^2} v_{2f}^2 + 2 \frac{m_2}{m_1} v_{1f} v_{2f} \left[\cos \theta_1 \cos \theta_2 + \sin \theta_1 \sin \theta_2 \right]$
 $\cos(\theta_1 - \theta_2)$

$\boxed{v_{1i}^2} = \boxed{v_{1f}^2} + \frac{m_2^2}{m_1^2} v_{2f}^2 + 2 \frac{m_2}{m_1} v_{1f} v_{2f} \cos(\theta_1 - \theta_2)$ (a)

Direct consequence of $\vec{P}_i = \vec{P}_f$!

Divide both side of eq ③ by $\frac{1}{2} m_1$:

$\boxed{v_{1i}^2} = \boxed{v_{1f}^2} + \frac{m_2}{m_1} v_{2f}^2$ (b)
 $KE_i = KE_f$

Similar terms in (a) & (b) \rightarrow (a) - (b)!

(a)-(b):
$$0 = \left(\frac{m_2}{m_1} - \frac{m_2}{m_1} \right) \sqrt{v_{2f}^2} + 2 \sqrt{\frac{m_2}{m_1}} v_{1f} \sqrt{v_{2f}^2} \cos(\theta_1 - \theta_2)$$

$$\sqrt{\frac{m_2}{m_1}} \left(\frac{m_2}{m_1} - 1 \right)$$

(a)-(b) :
$$0 = \left(\frac{m_2}{m_1} - 1 \right) v_{2f} + 2 v_{1f} \cos(\theta_1 - \theta_2)$$

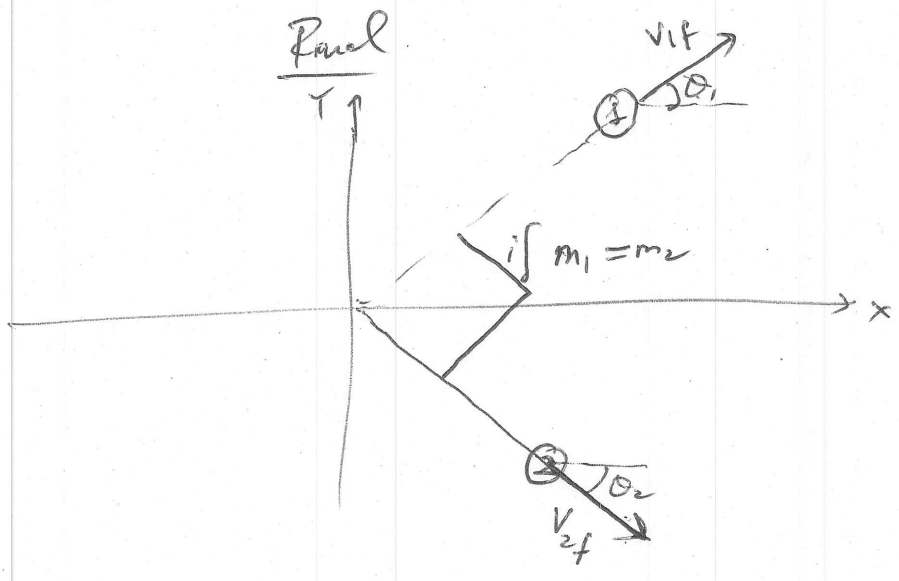
2D Elastic Collision $\begin{cases} \vec{P}_i = \vec{P}_f \\ KE_i = KE_f \end{cases}$

If $m_2 = m_1$

$$\rightarrow 2 v_{1f} \cos(\theta_1 - \theta_2) = 0 \rightarrow \cos(\theta_1 - \theta_2) = 0$$

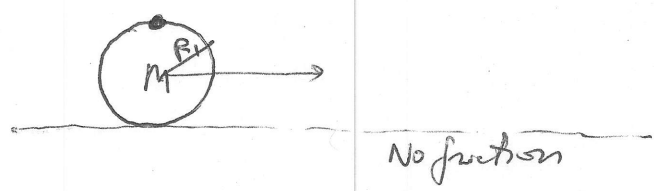
$$\Rightarrow \theta_1 - \theta_2 = 90^\circ$$

Direction of the objects after collision form 90° angle if $m_1 = m_2$!



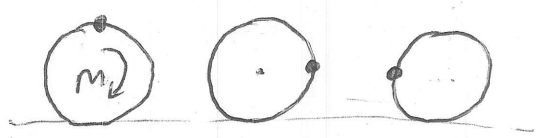
Ch. 10 Rotational Motion

Translation

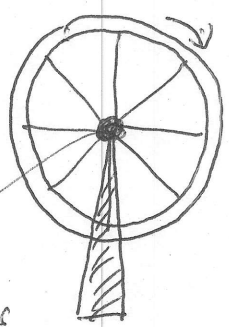


- Sliding balls of equal masses M but different radii $R_2 < R_1$;
- Same translational motion: since their motions (translations) can be described as that of a point-like particle of mass M located @ their center of masses (CM)
- Sliding or translation: top point (as shown) always stay @ top.

Rotation

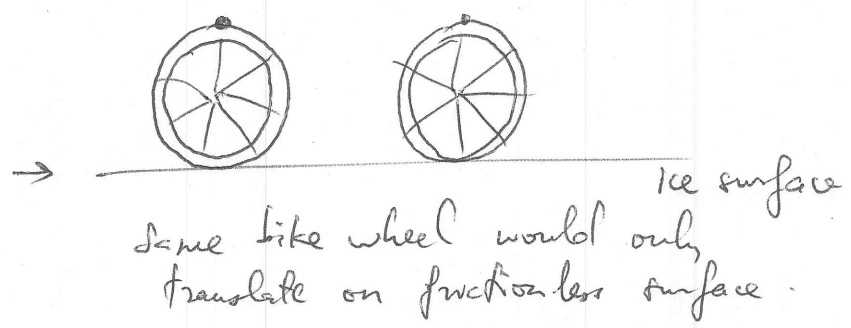


- Pure rotation, because of friction we see rolling balls. Rolling motion involves both rotation & translation.
- ① When we look at rotation size does matter.
- ② Are translation. Are rotations always associated with translations? No!



Center or axis of rotation or pivot point

→ A bike wheel on a support rotates without any translation!



Other situations:

- 1) Car wheels under normal ~~static~~ road conditions:
 - Rolling motion: both translation & rotation
- 2) Car wheels stuck in sand:
 - Only rotation
- 3) Car wheels while brake pedal is applied (no ABS): only translation
 - " " " " " " " (ABS): both rotation & translation.

ABS: anti-blocking braking system: advantages:

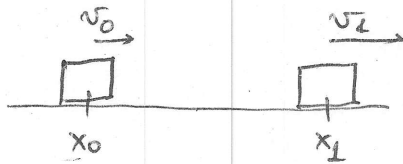
w/o ABS: can only rely on friction to lose final KE to come to a stop.

w/ ABS: in addition to friction can lose some of the final KE into rotational motion of wheels

- stopping distance is shorter.
- more vehicle control at curves.

Translational Motion

(Change of position)



$$\bar{v} = \frac{v_0 + v_1}{2}$$

$$\bar{v} = \frac{\Delta x}{\Delta t} \quad (\text{average vel.})$$

Change of position over change of time ($\frac{m}{s}$)

$$v = \frac{dx}{dt} \quad (\text{instantaneous vel.})$$

$$v = v_0 + a \cdot t \quad (1)$$

a : linear acceleration ($\frac{m}{s^2}$)

$$x = x_0 + v_0 \cdot t + \frac{1}{2} a t^2 \quad (2)$$

$$\frac{v^2 - v_0^2}{x - x_0} = 2 \cdot a \quad (3)$$

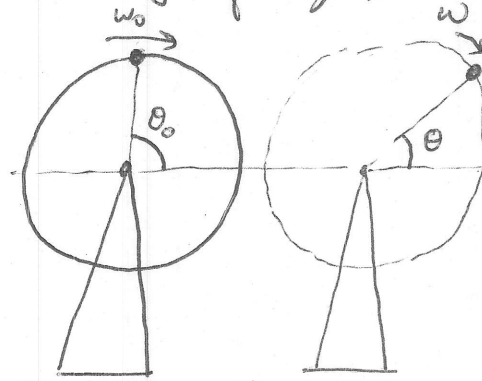
$$F_{net} = m \cdot a$$

↓ inertia

(constant mass)

Rotational Motion

(change of angle)



$$\bar{\omega} = \frac{\omega_0 + \omega}{2}$$

Same system (it did not translate!)

ω : angular velocity

$$\bar{\omega} = \frac{\Delta \theta}{\Delta t} \quad (\text{change of angle over change of time})$$

$\frac{\text{radian}}{s}$ or $\frac{\text{rad}}{s}$

(2π radians in 360°)

$$\omega = \frac{d\theta}{dt} \quad (\text{instantaneous ang. velocity})$$

$$\omega = \omega_0 + \alpha \cdot t \quad (1)$$

α : angular acceleration ($\frac{\text{rad}}{s^2}$)

$$\theta = \theta_0 + \omega_0 \cdot t + \frac{1}{2} \alpha t^2 \quad (2)$$

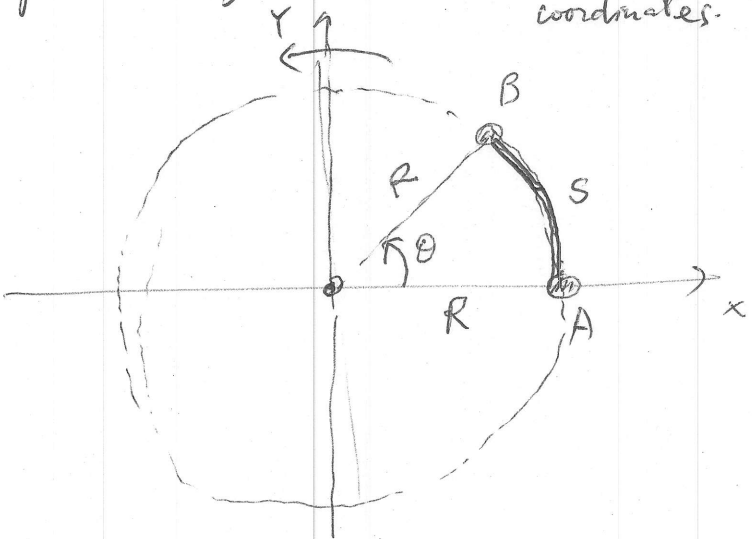
$$\frac{\omega^2 - \omega_0^2}{\theta - \theta_0} = 2\alpha \quad (3)$$

$$\tau_{net} = I \cdot \alpha$$

τ_{net} = "tau" = net torque } size of object
 I = moment of inertia } does matter

- Angular velocity (ω) and acceleration (α)
- Net torque & moment of inertia (size matters)

1) Any connection b/w angular velocity & linear velocity of a point on object?
 Point A on a rotating disk w/ origin of coordinates.



- a) From A to B angle went from 0 to θ (ccw)
- b) A has changed position along a circular trajectory of radius R, an arc length s

b) $\theta = \frac{s}{R}$ \rightarrow linear motion on circular trajectory
 \downarrow
 angular rotation

$$\frac{d}{dt} \left[\theta = \frac{s}{R} \right] \rightarrow \frac{d\theta}{dt} = \frac{1}{R} \frac{ds}{dt} \rightarrow \boxed{\omega = \frac{v}{R}}$$

For a point on the rotating disk.

$$\omega = \frac{v}{R} = \frac{\frac{m}{s}}{m} = \frac{1}{s}$$

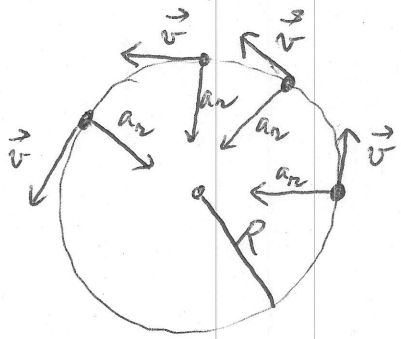
\downarrow
rpm (rev per minute)

\rightarrow correct since angle is dimensionless ($[\theta] = 1$)

2) Angular acceleration: α ("alpha")

$\bar{\alpha} = \frac{\Delta \omega}{\Delta t}$, $\alpha = \frac{d\omega}{dt}$ ($\frac{\text{rad}}{\text{s}^2} = \frac{1}{\text{s}^2} = \text{s}^{-2}$)

→ UCM (Uniform circular motion)
 - constant speed (tangential)
 - constant radial acceleration a_r (to change dir of motion)
 - tangential acceleration is 0



$a_r = \frac{v^2}{R}$ (UCM: $v = \text{constant}$, $R = \text{constant}$)
 $a_t = 0$
 $\Rightarrow a_r = \frac{\omega^2 R^2}{R} = \omega^2 R$
 $\omega = \frac{v}{R}$

What is α ?
 $\alpha = \frac{d\omega}{dt} = \frac{1}{R} \frac{dv}{dt} = \frac{1}{R} a_t = 0$

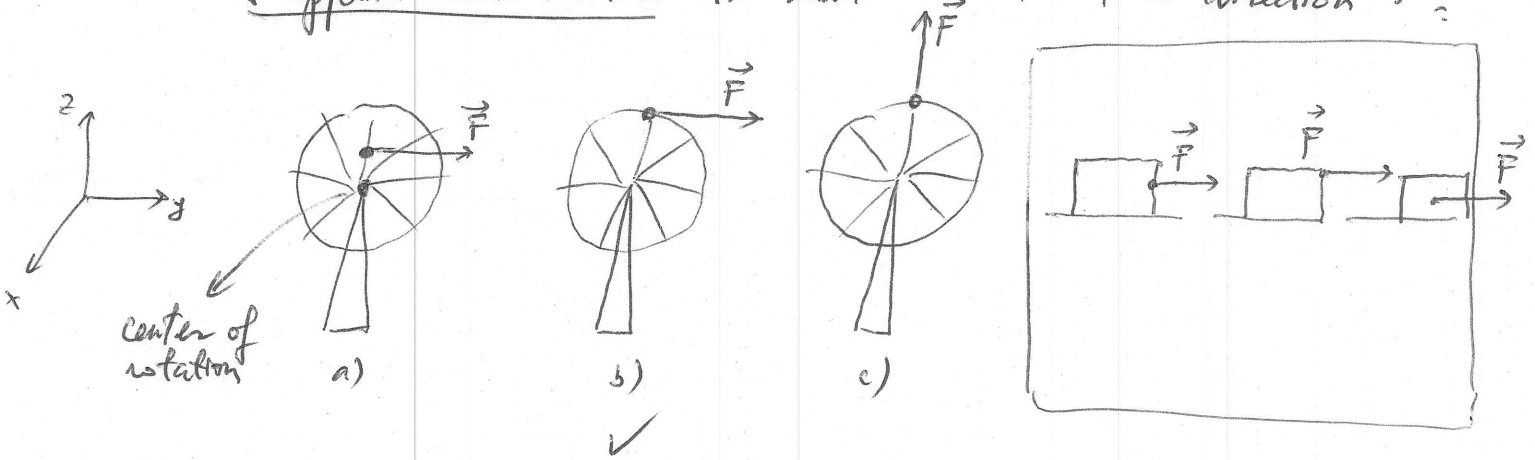
Makes sense: in UCM = constant $v \rightarrow$ constant ω ($\omega = \frac{v}{R}$)
 $\rightarrow \alpha = 0$

→ Non-UCM or non-uniform circular motion:

$a_r = \text{still there (whenever there is a circular motion!)} = R\omega^2$
 $a_t = \frac{dv}{dt} = \frac{d(\omega R)}{dt} = R \frac{d\omega}{dt} = R\alpha$

3) Torque $\vec{\tau}$ (tau) :

Is needed in rotations b/c size matters: → where a force is applied does matter in addition to the force direction !



$$\vec{\tau} = \vec{r} \times \vec{F}$$

\vec{r} : position vector of the force application point, w.r.t pivot or center of rotation.
 \times : "cross product": product of two vectors that is another vector that is perpendicular to both

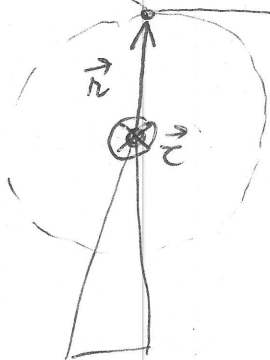
$$= r F \sin \theta \hat{e}$$

θ : angle b/w \vec{r} & \vec{F}
 \hat{e} : unit vector that is perpendicular to the plane formed by \vec{r} & \vec{F} direction given by Right Hand Rule (RHR)

Unit: Nm

Torque is a vector whose magnitude is the product of the magnitudes of the position vector & force vector times the sine of the angle b/w them. Direction of the torque is perpendicular to both \vec{r} & \vec{F} and given by RHR.

Cross-product & RHR
Force application point

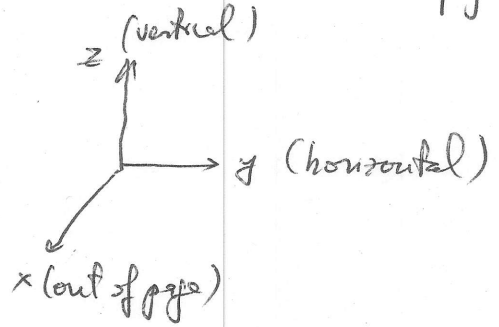


$$\vec{\tau} = \vec{r} \times \vec{F}$$

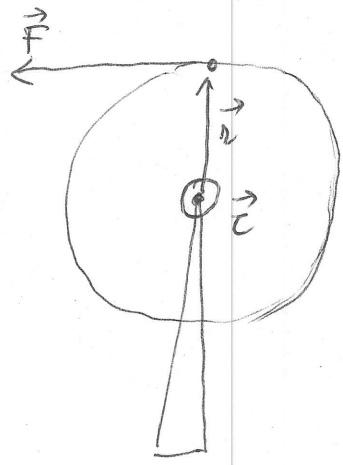
As shown: \vec{r} & \vec{F} are in plane of this page. $\rightarrow \vec{\tau}$ is perpendicular to the page.

RHR: align RH ~~thumb~~ ^{fingers} along 1st vector (\vec{r}), as you turn close RH fingers toward direction of 2nd vector (\vec{F}), RH thumb points in direction of torque:

As shown: $\vec{\tau}$ = into page \otimes



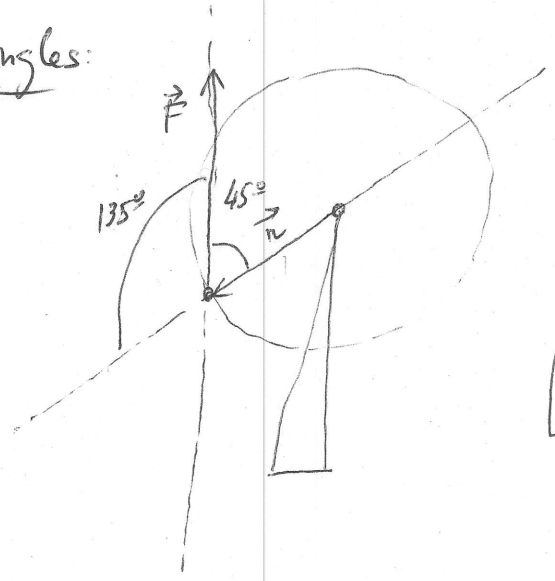
$$\vec{\tau} = rF \sin 90^\circ (-\hat{i}) = -rF\hat{i}$$



RHR: $\vec{\tau}$ = out of page \odot

$$\vec{\tau} = +rF\hat{i}$$

Angles:



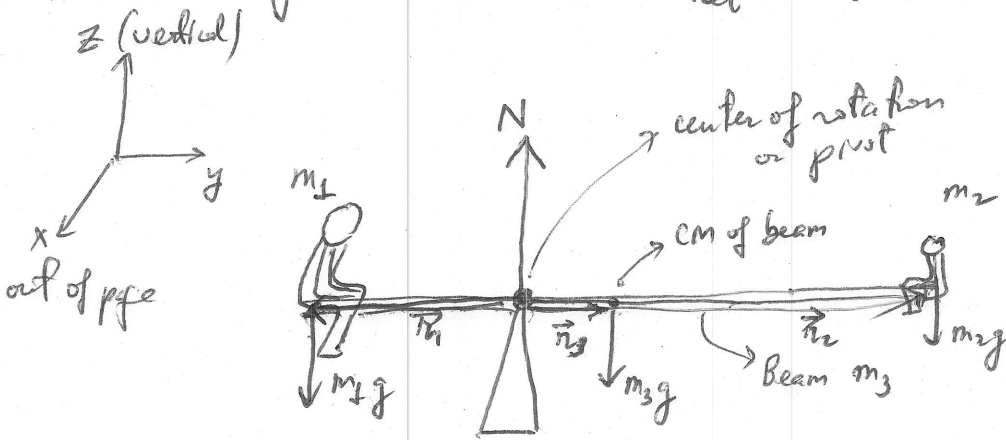
$$\tau = rF \sin \theta \quad (\text{Magnitude: } +)$$

$$\sin 45^\circ = - \sin 135^\circ$$

$$\tau = rF |\sin \theta|$$

Balancing a beam:

$$\vec{\tau}_{net} = 0 = I \cdot \alpha \quad (\alpha = 0)$$



Forces applying on beam:

- its weight m_3g
- m_1g
- m_2g
- N : normal force from pivot point.

Torques apply on beam:

- $\vec{\tau}_1 = r_1 m_1 g \frac{\sin 90}{1} \hat{i}$
- $\vec{\tau}_2 = r_2 m_2 g \frac{\sin 90}{1} (-\hat{i})$
- $\vec{\tau}_3 = r_3 m_3 g \frac{1}{1} (-\hat{i})$
- No torque by N ($\vec{r} = 0$)

\hookrightarrow wrt to center of rotation or pivot

$$\vec{\tau}_{net} = \vec{\tau}_1 + \vec{\tau}_2 + \vec{\tau}_3 = (r_1 m_1 g - r_2 m_2 g - r_3 m_3 g) \hat{i} = 0$$

Example:

- beam length $L = 3m$
- pivot at $0.75m$ from left edge
- $m_1 = 75 \text{ kg}$; $m_2 = 25 \text{ kg}$

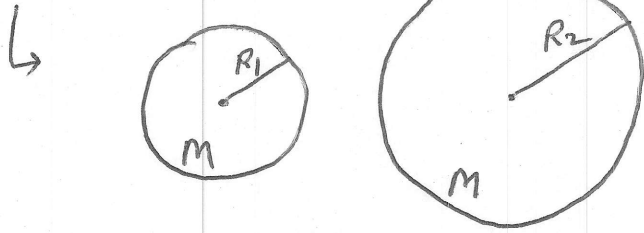
$$m_3 = \frac{(r_1 m_1 - r_2 m_2) g}{r_3 g} = \frac{0.75 \cdot 75 - 2.25 \cdot 25}{0.75} = 25 \left(\frac{2}{2} - \frac{3}{2} \right) = 0 \text{ kg.}$$

$$r_1 = 0.75m, \quad r_2 = 2.25m, \quad r_3 = 0.75m$$

4) Moment of Inertia (I)

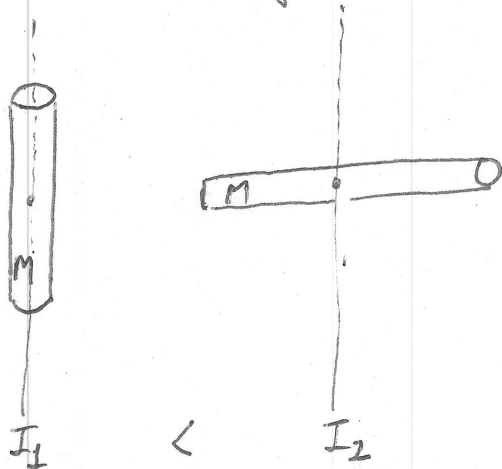
$(\checkmark \tau_{net} = I \cdot \checkmark \alpha)$

(size of mass distribution matters!)



$I_1 < I_2$

location of the axis of rotation also matters:



$I_1 < I_2$

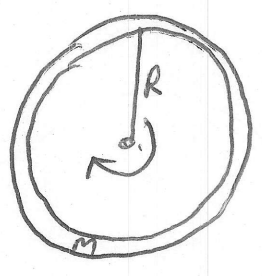
$I = \begin{cases} \text{Discrete system: } \sum_i m_i r_i^2 \\ \text{Continuous system: } \int dm r^2 \end{cases}$

Moment of inertia for round & symmetrical objects:
disks, rods, spheres, etc...

$I = \alpha MR^2$ $\left\{ \begin{array}{l} M: \text{total mass of object} \\ R: \text{radius of mass distribution wrt center of rotation} \end{array} \right.$

- ↳ 1) Sphere wrt center axis $\rightarrow \alpha = \frac{2}{5}$ (proved by $\int dm r^2$)
- 2) Cylinder wrt center axis $\rightarrow \alpha = \frac{1}{2}$
- 3) Thin rod of length L wrt middle axis $\alpha = \frac{1}{12}$; $R = L/2$ or $I = \frac{1}{12} ML^2$

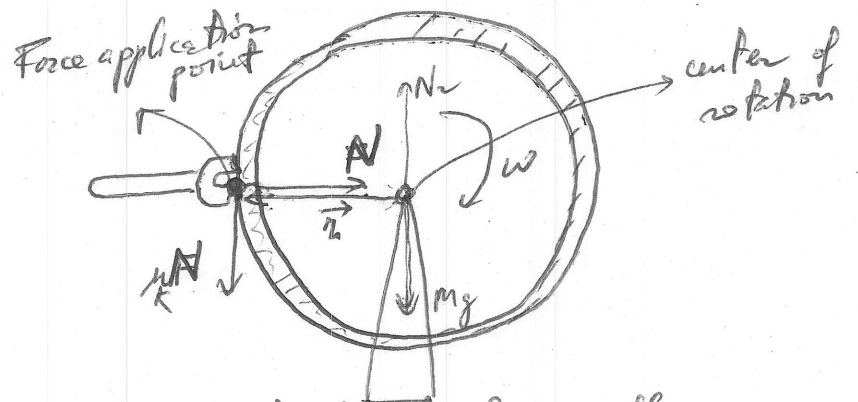
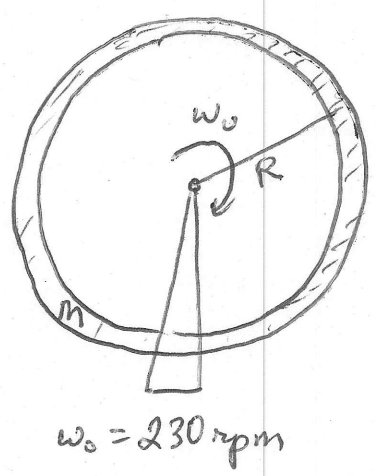
4) Ring of mass M & radius R w/ center axis.



$\alpha = 1$
 or $I = MR^2$

Example:
 10.58

Spinning bike wheel: $R = 0.33m, M = 1.9kg \rightarrow I = MR^2$



gets slowed down with
 a wrench at normal direction
 $N = 2.7N$ during $\Delta t = 3.1s$
 $\mu_k = 0.46$
 (b/w wrench & tire)

Solve for final angular speed ω (after $N = 2.7N$ is applied during $\Delta t = 3.1s$):

$$\omega = \omega_0 - \alpha \Delta t$$
 (deceleration happens while wrench is applied)

As 2nd Newton's law was center piece to solve linear motion, its analog $\tau_{net} = I\alpha$ is the center piece to solve rotational motion.

$$\tau_{net} = I \cdot \alpha \rightarrow \alpha = \frac{\tau_{net}}{I} = \frac{\mu_k N R}{MR^2}$$

τ_{net} = sum of torques on bike wheel:

→ There are 4 forces on bike wheel: $N, \mu_k N, Mg, N_2$
 → There is 1 torque (wrt center rotation): $\tau_{net} = R \mu_k N \sin 90^\circ$
weight. Normal @ pivot

N_2 & M_2 apply no torque since center of rotation is also force application point ($\vec{r} = 0!$)

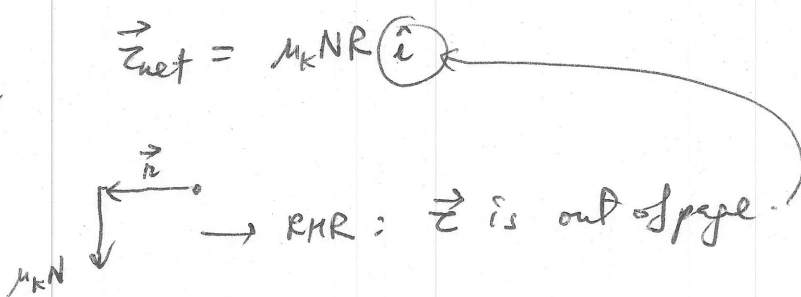
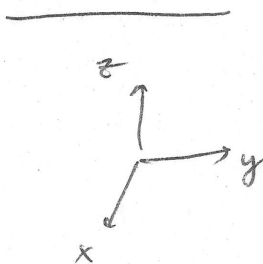
N applies no torque: although center of rotation is R from force application point $\theta = 180^\circ \Rightarrow \sin \theta = 0$

$$\omega = \omega_0 - \frac{\tau_{net}}{I} \cdot \Delta t = \omega_0 - \underbrace{\frac{\mu_k N R}{MR^2}}_{\Delta \omega} \Delta t$$

$$\Delta \omega = \frac{0.46 \cdot 2.7 \cdot 0.33}{1.9 \cdot 0.33^2} \cdot 3.1 = 6.14 \frac{\text{rad}}{s} \cdot \frac{1 \text{ rev}}{2\pi \text{ rad}} \cdot \frac{60 \text{ s}}{1 \text{ min}}$$

SI units. $= 58.6 \frac{\text{rev}}{\text{min}} = 58.6 \text{ rpm}$

$$\rightarrow \omega = \omega_0 - \Delta \omega = 230 - 58.6 = 171 \text{ rpm.}$$

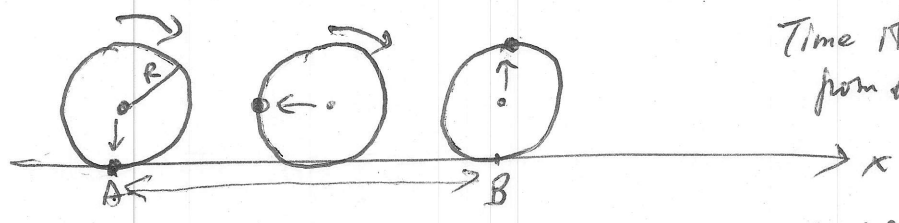


$$\vec{\tau}_{net} = \mu_k N R \hat{z}$$

RHR: \hat{z} is out of page.

Rolling Motion:

↳ Non-skidding



Time it takes to go from A to B is Δt .

Disk rolling on horizontal surface, non-skidding

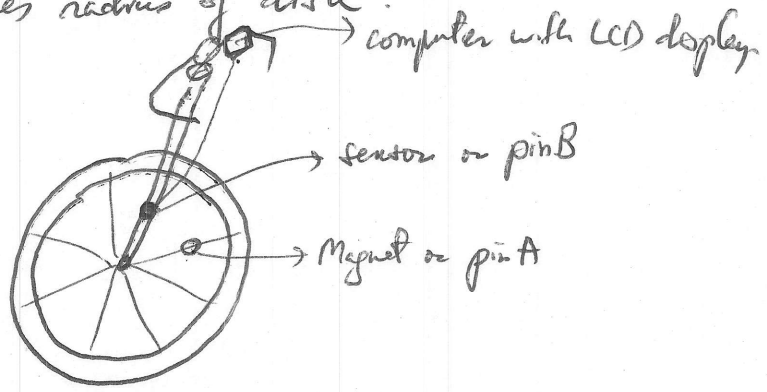
- 1) Center of Mass of disk has traveled from A to B
- 2) Disk in contact with surface at all time (non-skidding)
 - ↳ half of its circumference has been in contact with the surface = $\Delta x_{AB} = \pi R$

$$v_{cm} = \frac{\Delta x_{AB}}{\Delta t} = \frac{\pi R}{\Delta t} = \omega R \rightarrow \boxed{v_{cm} = \omega R}$$

angular speed of disk = $\omega = \frac{\Delta \theta}{\Delta t} = \frac{\pi}{\Delta t}$

In rolling motion: linear speed of CM is angular speed times radius of disk.

Apply this to a ring:
 ↳ "bike speedometer"



Every time pin A passes pin B: computer assumes wheel has completed one revolution:
 → calculates ω in rpm, then use rolling motion formula $v_{cm} = \omega \cdot R$ to display linear speed.

Kinetic energy in rolling motion:

① Linear motion

$$\frac{1}{2}mv^2$$



Non-friction
Disk skidding w/o rotation on a frictionless surface

② Rotational Motion

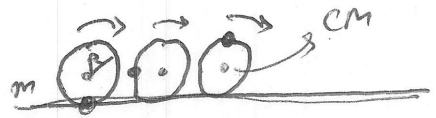
$$\frac{1}{2}I\omega^2$$



Disk is rotating freely on a support (no translational motion)

③ Rolling motion

$$\frac{1}{2}mv^2 + \frac{1}{2}I\omega^2$$



Disk rolling on surface with friction: translation of CM + rotation wrt CM

KE for rolling motion:

$$v_{CM} = \omega \cdot R \rightarrow \omega = \frac{v_{CM}}{R}$$

$$KE = \frac{1}{2}mv_{CM}^2 + \frac{1}{2}I\left(\frac{v_{CM}}{R}\right)^2 = \frac{1}{2}\left(m + \frac{I}{R^2}\right)v_{CM}^2$$

For a disk = $I = \frac{1}{2}MR^2$

$$m + \frac{I}{R^2} = m + \frac{\frac{1}{2}mR^2}{R^2} = m + \frac{m}{2} = \frac{3}{2}m$$

Disk: Sliding
wrt center axis
 $KE = \frac{1}{2}mv_{CM}^2$

Rolling
 $KE = \frac{1}{2}\left(\frac{3}{2}m\right)v_{CM}^2$

ABS: anti-locking brakes

In term of KE rolling increases mass ~~total~~ inertia a factor of 1.5 (150%)
 { w/o ABS: inertia mass is m
 w/ ABS: inertia mass is 1.5m per wheel.
 → shorter stopping distance!

Ch 11 Rotational Vectors & Angular Momentum:

Linear Motion

\vec{p} = linear momentum
 $\vec{p} = m\vec{v}$

2nd Newton's Law:

$$\vec{F}_{net} = \frac{d\vec{p}}{dt}$$

$$\vec{F}_{net} = 0 \Rightarrow \vec{p}_i = \vec{p}_f$$

Total linear momentum is conserved if $\vec{F}_{net} = 0$

↓
Collisions!

Rotational Motion

\vec{L} = angular momentum

$$\vec{L} = \vec{r} \times \vec{p}$$

→ cross product b/w pos. vector (wrt center of rotation) & linear momentum

→ as with the torque, \vec{L} is perpendicular to both \vec{r} & \vec{p} with direction given by RHR

Analogy of 2nd Newton's Law:

$$\vec{\tau}_{net} = \frac{d\vec{L}}{dt}$$

$$\vec{\tau}_{net} = 0 \Rightarrow \vec{L}_i = \vec{L}_f$$

Total angular momentum is conserved if $\vec{\tau}_{net} = 0$

And $\vec{\tau}_{net} = I \cdot \vec{\alpha}$
 $\hookrightarrow \boxed{\vec{L} = I \cdot \vec{\omega}}$

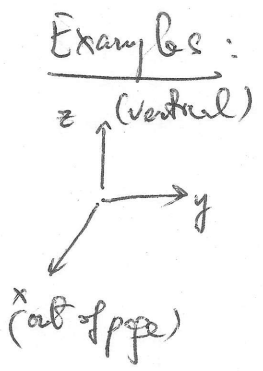
Two cross products: with a defined center of rotation

$$\vec{\tau} = \vec{r} \times \vec{F}$$

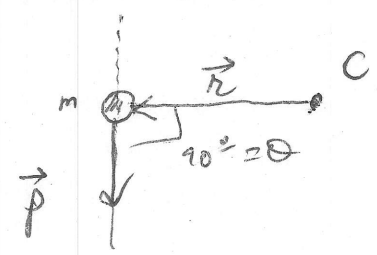
$$\vec{L} = \vec{r} \times \vec{p}$$

\vec{r} = position vector, from center of rotation to force application point

\vec{r} = position vector, from center of rotation to the position of mass m that has linear momentum \vec{p}

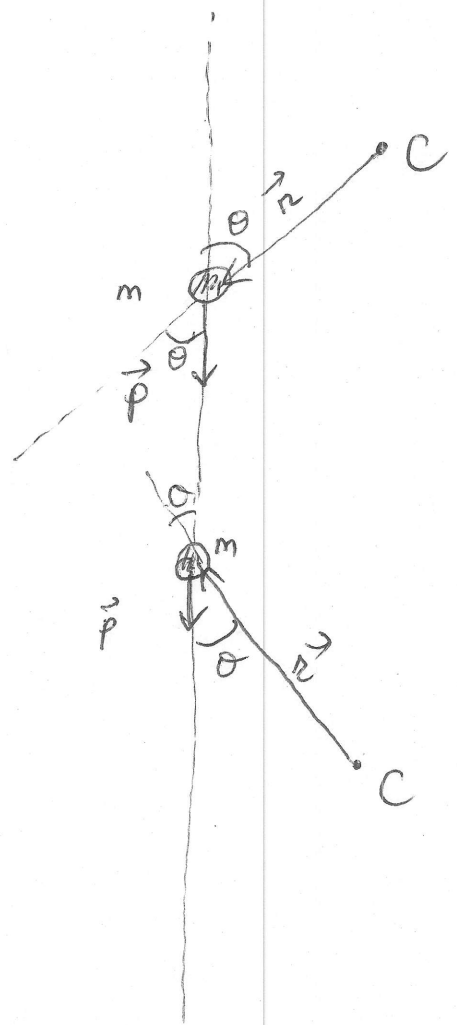


Rotation with centers C:
 object of mass m moving along $-z$ ($-k$)
 (an object does not need to rotate to have an angular momentum)



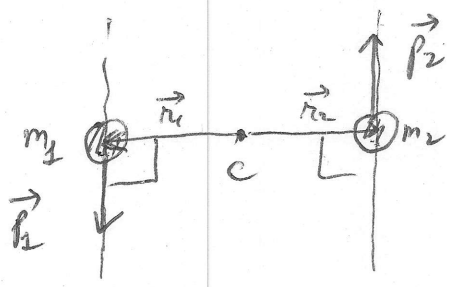
Angular momentum for this object: $\vec{L} = \vec{r} \times \vec{p} = rp \hat{i}$
 given by RHR

What is the angular momentum now?



$\vec{L} = \vec{r} \times \vec{p} = rp \sin \theta \hat{i}$
 by RHR

$\vec{L} = rp \sin \theta \hat{i}$
 by RHR

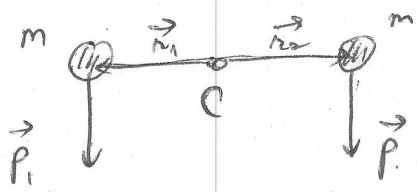


System with 2 masses →
 Total angular momentum \vec{L} is

$$\begin{aligned} \vec{L} &= \vec{L}_1 + \vec{L}_2 = \vec{r}_1 \times \vec{p}_1 + \vec{r}_2 \times \vec{p}_2 \\ &= r_1 p_1 \hat{i} + r_2 p_2 \hat{i} \\ &= 2rp \hat{i} \end{aligned}$$

↓
 C right of both masses $r_1 = r_2 = r$
 both masses with same speed $\vec{p}_1 = \vec{p}_2 = \vec{p}$

In a car: this happens in the steering wheel. ~~when~~
 it turns → makes sense there is a total \vec{L} !



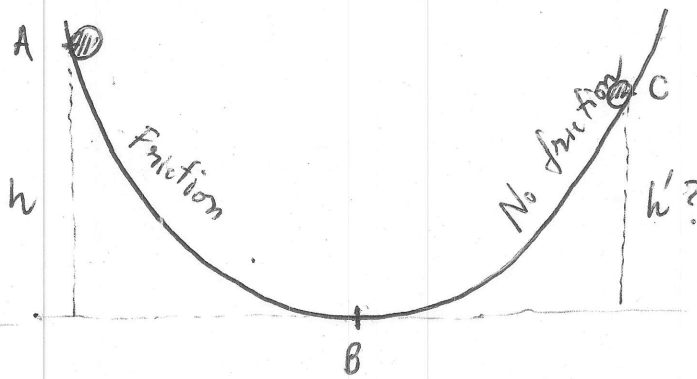
$$\vec{L} = \cancel{rp} rp \hat{i} - rp \hat{i} = 0$$

↓
 makes sense!

10.64

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Solid ball \rightarrow sphere $I = \frac{2}{5}MR^2$
 M, R



Point C:
 Above A
 Same height as A
 Below A

Related \rightarrow Energy lost due to friction.
 - There was rotation besides translation b/w A & B.

Find h' :

AB: Rolling motion
 (Translation & rotation and $v_{cm} = \omega R$)

BC: Sliding motion
 (only translation)

Conservation of energy:

AB: $Mgh = \frac{1}{2}Mv_B^2 + \frac{1}{2}I\omega^2$ (1)

BC: $\frac{1}{2}Mv_B^2 = Mgh'$ (2)

$h' = \frac{v_B^2}{2g}$

Solve for v_B in (1):

- (a) spherical solid ball rotating wrt its center axis $I = \frac{2}{5}MR^2$
- (b) Rolling motion b/w A & B: $\omega = \frac{v_B}{R}$

(1) $Mgh = \frac{1}{2}Mv_B^2 + \frac{1}{2} \cdot \frac{2}{5}MR^2 \cdot \frac{v_B^2}{R^2}$

$gh = \frac{v_B^2}{2} + \frac{v_B^2}{5} = \frac{v_B^2}{2} \left(1 + \frac{2}{5}\right) \rightarrow v_B^2 = \frac{10}{7}gh$

(2) $h' = \frac{\frac{10}{7}gh}{2g} = \frac{5}{7}h < h$

Note: this confirms our prediction that $h' < h$ since some initial gravitational potential energy was lost due to friction (which went into the rotational part of rolling motion b/w A & B)

10.65

Back ground information: Parallel Axis Theorem for Moments of Inertia

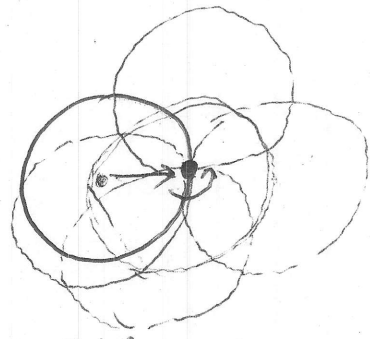
Rotation of a ~~sphere~~ **disk** as view from above:



Rotates w/ its center

axis:

$$I_1 = \frac{1}{2} MR^2$$



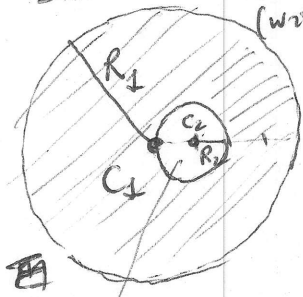
Rotates w/ an axis on the outer edge

$$I_2 = \frac{1}{2} MR^2 + MR^2 = \frac{3}{2} MR^2$$

Parallel Axis Theorem: when the axis is displaced by a distance R : the moment of inertia w/ the new axis is increased by $M \cdot R^2$

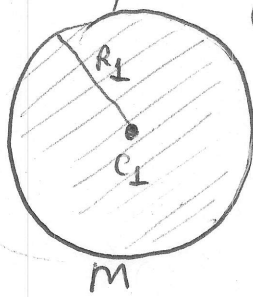
\downarrow mass of object \downarrow displacement squared.

Disk with hole: I (w/ C_1)



hole of radius $R_2 = \frac{R_1}{4}$

Disk w/ no hole: I_1 (w/ C_1)



Smaller disk: I_2 (w/ C_1)



- Mass of disk before hole was drilled is M
- C_1 = center of disk and center of rotation
- C_2 = center of the hole

$$I = I_1 - I_2$$

$$= \frac{1}{2} MR_1^2 - \left(\frac{1}{2} m R_2^2 + m R_2^2 \right)$$

Parallel axis Theorem

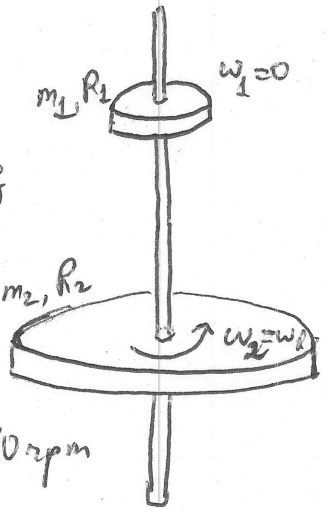
Note: $\left\{ \begin{array}{l} R_2 = \frac{R_1}{4} = \frac{R}{4} \\ m = \frac{\pi R_2^2}{\pi R_1^2} M = \frac{1}{16} M \end{array} \right.$

$$\rightarrow I = \frac{1}{2} MR_1^2 - \frac{3}{2} m R_2^2 = \frac{1}{2} MR_1^2 - \frac{3}{2} \frac{M}{16} \frac{R_1^2}{16} = \frac{1}{2} MR_1^2 \left(1 - \frac{3}{16^2} \right) = \frac{1}{2} MR_1^2 \left(\frac{253}{256} \right)$$

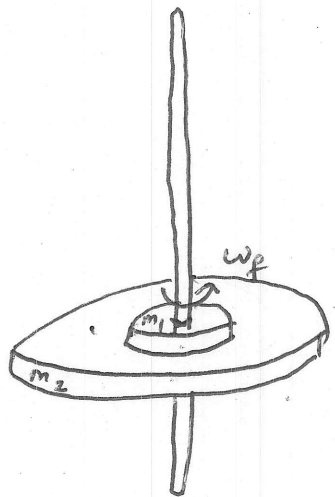
all these I 's (w/ to same central axis C_1)

11.51

- $m_2 = 0.44 \text{ kg}$
- $R_2 = 0.035 \text{ m}$
- $m_1 = 0.27 \text{ kg}$
- $R_1 = 0.023 \text{ m}$



Initial
(Only m_2 is rotating)
 $\omega_i = \omega_2$



Final
(both are rotating at same speed ω_f due to friction b/w the two disks)

Initial & final \rightarrow conservation of angular momentum (rotations are involved)

Check: System of disks (shaft's mass is ignored) w/o friction on shaft: $\vec{\tau}_{\text{net}} = 0 \rightarrow \boxed{\vec{L}_i = \vec{L}_f}$
"m1 drops freely"
Force applied in direction of rotation: none

a) $\omega_f?$

$$L_i = L_f$$

$$I_2 \omega_2 = I_1 \omega_f + I_2 \omega_f = (I_1 + I_2) \omega_f$$

$$\omega_f = \frac{I_2}{I_1 + I_2} \omega_2 = \frac{\frac{1}{2} m_2 R_2^2}{\frac{1}{2} m_1 R_1^2 + \frac{1}{2} m_2 R_2^2} \omega_2$$

$$\omega_f = \frac{0.44 \times 0.035^2}{0.44 \times 0.035^2 + 0.27 \times 0.023^2} 180 \text{ rpm}$$

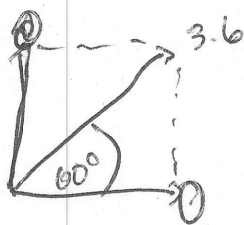
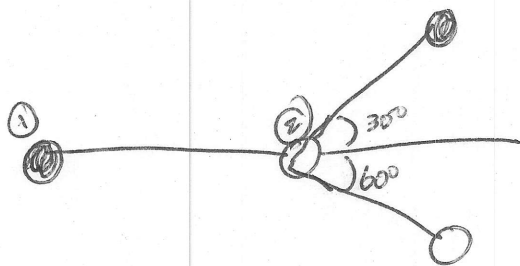
$$\omega_f = 142 \text{ rpm}$$

b) Fraction of energy lost due to friction:
 $KE_i > KE_f$ since some energy was lost to friction b/w the two disks! (for them to rotate at same final ω_f !)
Lost fraction = $\frac{KE_i - KE_f}{KE_i} = 1 - \frac{KE_f}{KE_i} = 1 - \frac{\frac{1}{2} (I_1 + I_2) \omega_f^2}{\frac{1}{2} I_2 \omega_2^2}$

$$\text{lost fraction} = 1 - \frac{\frac{1}{2}(m_1 R_1^2 + m_2 R_2^2) \omega_f^2}{\frac{1}{2} m_2 R_2^2 \omega_2^2}$$

$$= 1 - \frac{0.44 \cdot 0.035^2 + 0.27 \cdot 0.023^2}{0.44 \cdot 0.035^2} \cdot \frac{142^2}{180^2} = 0.210 \text{ or } \boxed{21\%}$$

conversion of rpm to $\frac{\text{rad}}{\text{s}}$
 is a scaling factor common
 to both numerator &
 denominator and would
 cancel out.



$$v_1 = \sin 60^\circ \cdot 3.6 = 3.12 \text{ m/s}$$

$$v_2 = \cos 60^\circ \cdot 3.6 = 1.8 \text{ m/s}$$

(10.37)

Translational & rotation KE



$$m = 0.15 \text{ kg}$$

$$R = 0.037 \text{ m}$$

$$v = 33 \frac{\text{m}}{\text{s}}$$

$$\omega = 42 \frac{\text{rad}}{\text{s}}$$

$$\text{Baseball} \rightarrow \text{sphere: } I = \frac{2}{5} MR^2$$

Fraction of rotational in the total KE?

$$\frac{KE_{\text{rot.}}}{KE_{\text{trans}} + KE_{\text{rot}}} = \frac{\frac{1}{2} I \omega^2}{\frac{1}{2} M v^2 + \frac{1}{2} I \omega^2}$$

$$\rightarrow = \frac{\frac{1}{2} \frac{2}{5} M R^2 \omega^2}{\frac{1}{2} M v^2 + \frac{1}{2} \frac{2}{5} M R^2 \omega^2}$$

$$= \frac{\frac{2}{5} R^2 \omega^2}{v^2 + \frac{2}{5} R^2 \omega^2}$$

$$= \frac{\frac{2}{5} 0.037^2 \cdot 42^2}{33^2 + \frac{2}{5} 0.037^2 \cdot 42^2}$$

$$= 8.86 \cdot 10^{-4} \text{ or } 0.0886\%$$

(Very small)

Note:

$$\omega R = 42 \cdot 0.037 = 1.554 \frac{\text{m}}{\text{s}} \rightarrow \text{This is the linear speed of a point on the surface of baseball}$$

$$v = 33 \frac{\text{m}}{\text{s}} \rightarrow \text{It's not } v \text{ (velocity of cm of ball) because this is not a rolling motion!}$$

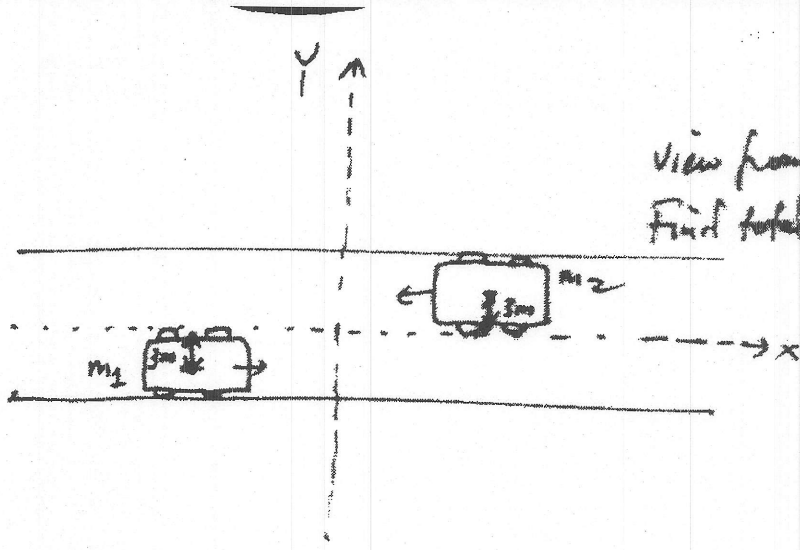
→ It's very small so most KE is in the translational motion

11.37

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view from above!
Find total $\vec{L} = \vec{L}_1 + \vec{L}_2$

$= \vec{r}_1 \times \vec{p}_1 + m_2 \vec{v}_2$



$m_1 = m_2 = 1800 \text{ kg}$

$90 \frac{\text{km}}{\text{h}} = 25 \frac{\text{m}}{\text{s}}$

$\vec{v}_1 = 25 \hat{i} \text{ m/s} ; \vec{r}_1 = x_1 \hat{i} - 3 \hat{j} \text{ (m)}$

$\vec{v}_2 = 25 (-\hat{i}) \text{ m/s} ; \vec{r}_2 = x_2 \hat{i} + 3 \hat{j} \text{ (m)}$

$$\vec{L} = \vec{L}_1 + \vec{L}_2 = \vec{r}_1 \times \vec{p}_1 + \vec{r}_2 \times \vec{p}_2$$

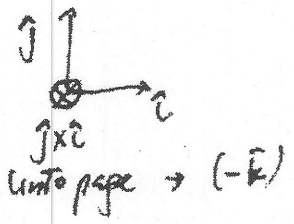
$$= (x_1 \hat{i} - 3 \hat{j}) \times \frac{25 \hat{i} \cdot 1800}{45000 \hat{i}} + (x_2 \hat{i} + 3 \hat{j}) \times \frac{25 (-\hat{i}) 1800}{-45000 \hat{i}}$$

$= 135000 \hat{k}$

$+ 135000 \hat{k} = 270000 \hat{k}$

$\hat{i} \times \hat{i} = 0$
 $\hat{j} \times \hat{i} = -\hat{k}$

$\frac{\text{kg} \cdot \text{m}^2}{\text{s}}$
 $= \text{J} \cdot \text{s}$



Equations:

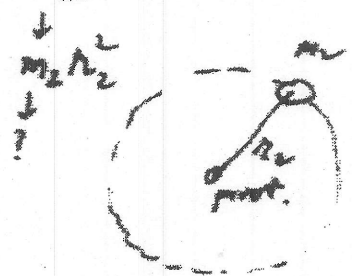
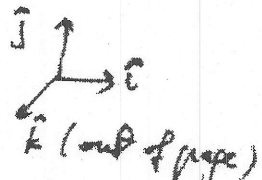
$$L_i = L_f$$

$$I_1 \omega_i + |\vec{r}_2 \times \vec{p}_2| = (I_1 + I_2) \omega_f$$

Clay

$$\vec{r}_2 \times \vec{p}_2 = (x_2 \hat{i} + y_2 \hat{j}) \times v_2 \hat{j} m_2$$

$$= x_2 v_2 m_2 \underbrace{\hat{i} \times \hat{j}}_{\hat{k}} = \frac{0.15 \times 1.3 \times m_2}{0.195} \hat{k}$$



$L_{table} = I_1 \omega_i (-\hat{k})$ opposite
 $L_{clay} = 0.195 m_2 (v_2 \hat{k})$ along
→ clay was opposing rotation!

Turn table

$I_1 \vec{\omega}_i = I_1 \omega_i (-\hat{k})$ (Direction by RHR: fingers turning $\rightarrow \vec{\omega}_i$; thumb is direction of $\vec{\omega}_i$)

$$\begin{aligned} -0.021 \times 0.29 + 0.195 m_2 &= - (0.021 + m_2 \cdot 0.15) \cdot 0.085 \\ + 0.021 \times 0.29 - 0.195 m_2 &= (0.021 + 0.0225) \cdot 0.085 - 0.0225 \\ + \frac{0.021 \times 0.29}{0.085} - 0.021 &= m_2 \left(0.0225 + \frac{0.195}{0.085} \right) \end{aligned}$$

$$m_2 = \frac{\frac{0.021 \times 0.29}{0.085} - 0.021}{0.0225 + \frac{0.195}{0.085}}$$

$$m_2 = 0.0218 \text{ kg} = 21.8 \text{ g}$$

→ Pay attention on the direction of angular momentum!

$$I = \frac{M R^2}{2} \Rightarrow M = \frac{2I}{R^2} = \frac{M}{16}$$