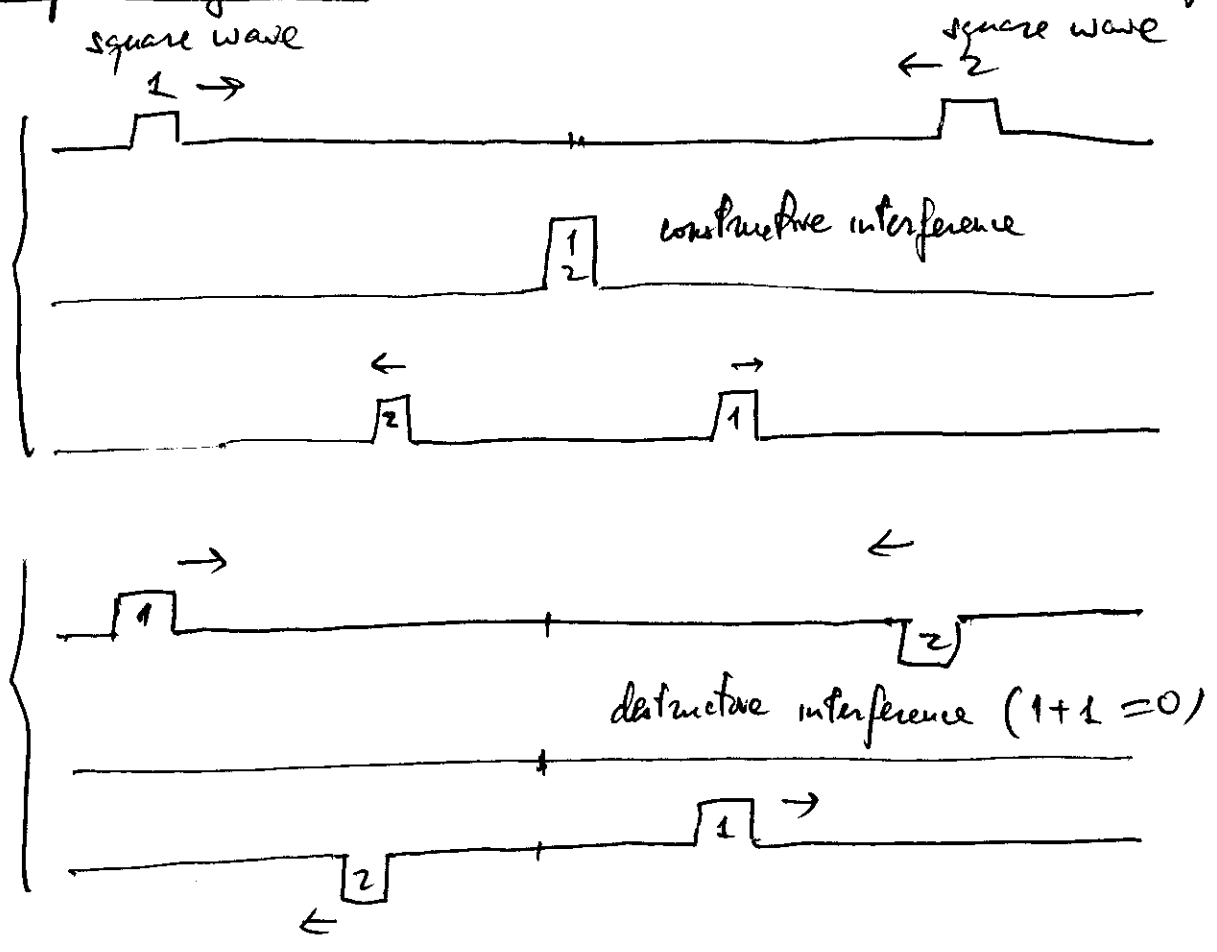


Ch 14 Wave Motion (Cont.)

Beat Phenomenon & Standing Waves & Doppler Effect

Superposition of Waves: constructive & destructive interference



Beat Phenomenon: (superposition of 2 sound waves) with slightly different frequencies

- Physical observation: a) Twin engine airplanes (propellers): should go at equal angular frequencies, however there is normally a small difference in frequencies. We can't hear the oscillations from either engine but we can hear the beating between the two.
- Math description
 - b) Old fashion way of tuning string instruments
 - ↳ Combine 2 transverse waves: $y(x,t) = A \sin(kx - \omega t)$

2 waves @ $x=0$ with frequencies ω_1 & ω_2

$$\begin{cases} y_1(0,t) = A \sin(-\omega_1 t) \\ y_2(0,t) = A \sin(-\omega_2 t) \end{cases}$$

For twin-engine & string tuning $x=0$ is location of our ear drum.

↓ Superposition of these 2 waves @ $x=0$:

$$y_1(0,t) + y_2(0,t) = -A [\sin \omega_1 t + \sin \omega_2 t]$$

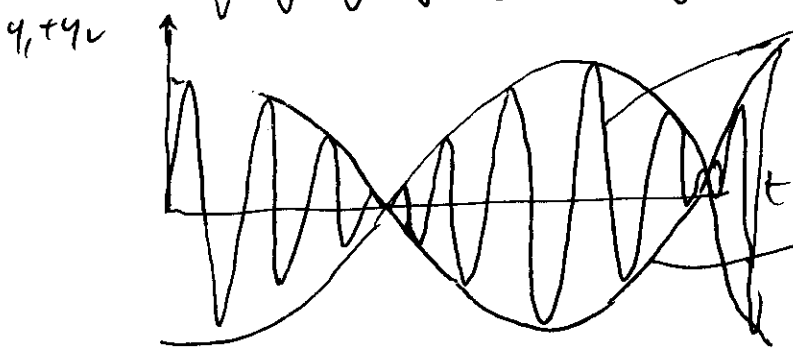
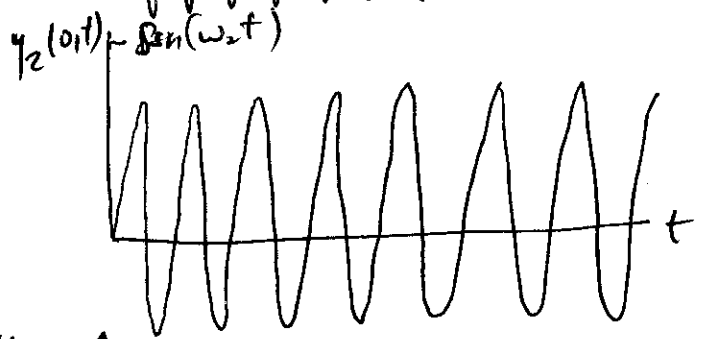
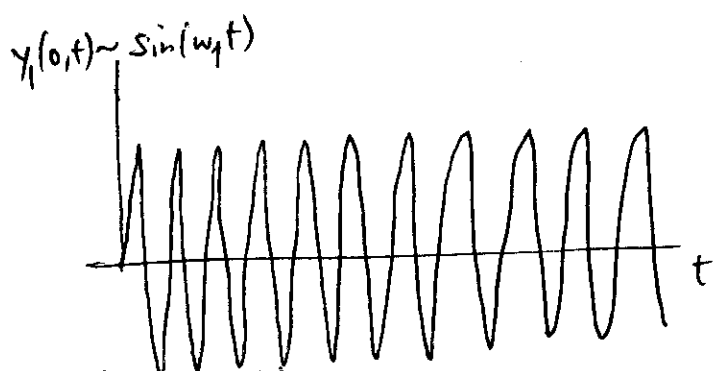
Trigonometry Identity: $\sin \alpha + \sin \beta = 2 \sin\left(\frac{\alpha+\beta}{2}\right) \cos\left(\frac{\alpha-\beta}{2}\right)$

$$= -2A \sin\left(\frac{\omega_1+\omega_2}{2}t\right) \cos\left(\frac{\omega_1-\omega_2}{2}t\right)$$

High frequency osc. low frequency modulation

When ω_1 & ω_2 are very similar

- 1) The average \approx either one
- 2) The difference is \ll either one



$\sin\left(\frac{\omega_1+\omega_2}{2}t\right) \rightarrow$ High frequency oscillations

$\cos\left(\frac{\omega_1-\omega_2}{2}t\right) \rightarrow$ low frequency envelope.
↓ Beats.

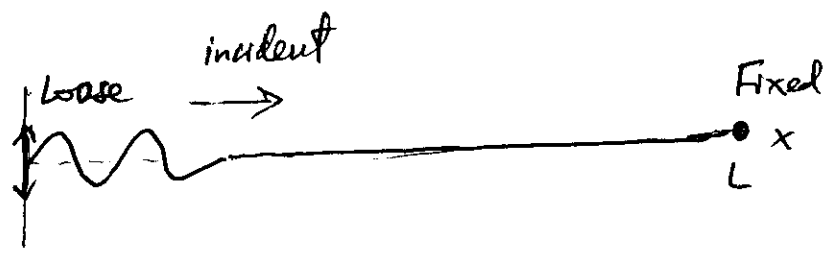
The high frequency oscillations are modulated by a very low frequency envelope which we can hear \rightarrow these are the beats.

Standing Waves:

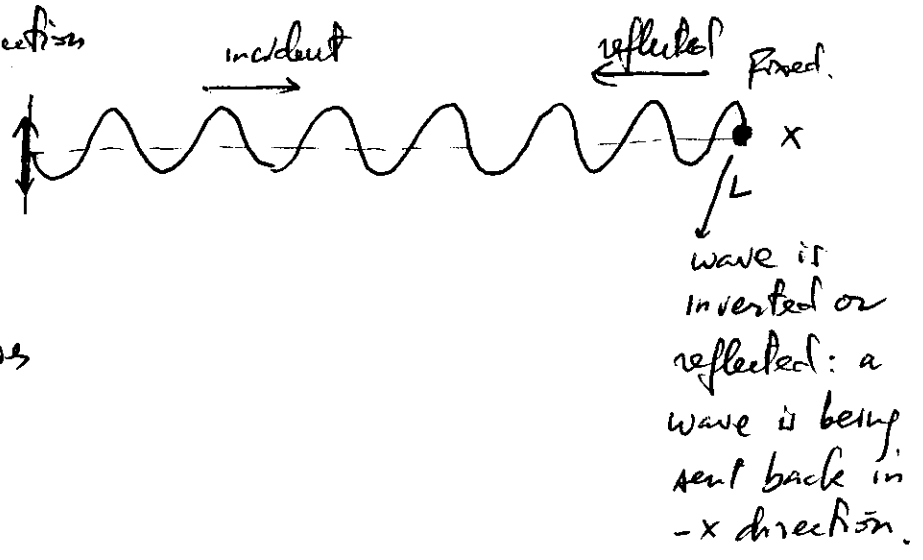
Physical observation:



- String along x-axis of length L. Fixed at x=L
- Transverse perturbation along y-axis @ loose end @ x=0 → transverse wave propagating in x-direction



- Reflected wave travels in -x direction
- Incident & reflected waves add up to standing waves not moving



Math description: $y_T(x,t) = y_{inc}(x,t) + y_{refl.}(x,t)$

$$= A \cos(kx - \omega t) \quad - \quad A \cos(kx + \omega t)$$

propagation in +x propagation in -x
 Reflected wave @ fixed end x=L

Trigonometry identity: $\cos \alpha - \cos \beta = -2 \sin \frac{\alpha + \beta}{2} \sin \frac{\alpha - \beta}{2}$

$$y_T(x,t) = -2A \sin \left(\frac{kx - \omega t + kx + \omega t}{2} \right) \sin \left(\frac{kx - \omega t - kx - \omega t}{2} \right)$$

$$= -2A \sin kx \sin(-\omega t) = 2A \sin kx \sin \omega t$$

Note: there is no further $(kx - \omega t)$ in the argument of \sin or $\cos \rightarrow$ no propagation, only oscillation in space ($\sin kx$) & time ($\sin \omega t$). \hookrightarrow Standing waves.

A very important consequence from math description of standing waves:

The string is fixed @ $x=L \Rightarrow y_T(L,t) = 2A \sin kL \sin \omega t = 0$
@ all time t .

$\rightarrow \sin kL = 0 \iff kL = n\pi$ ($n=1,2,3,\dots$) (a multiple of π)

$\frac{2\pi}{\lambda} L = n\pi$

$\lambda = \frac{2L}{n}$ ($n=1,2,3,\dots$)

\rightarrow Only wave of certain wave lengths can stand in a string of length L

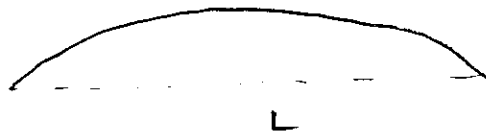
$\lambda_n = \frac{2L}{1}, \frac{2L}{2}, \frac{2L}{3}, \frac{2L}{4}, \dots$

\rightarrow For example, wave of wave length $\lambda = 2.5L$ can't stand in this string!

$\lambda_1 = 2L \rightarrow$ in one L we have half wave length

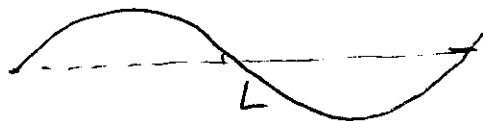
$\hookrightarrow \omega_1 = v \frac{\pi}{L}$

$\omega = kv \rightarrow \omega_1 = \frac{2\pi}{\lambda_1} v = \frac{2\pi}{2L} v$



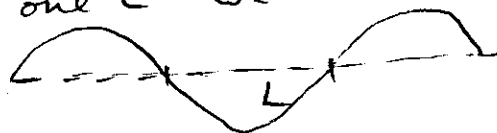
$\lambda_2 = L \rightarrow$ in one L we have a full wave length

$\hookrightarrow \omega_2 = 2\omega_1$



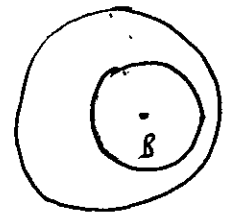
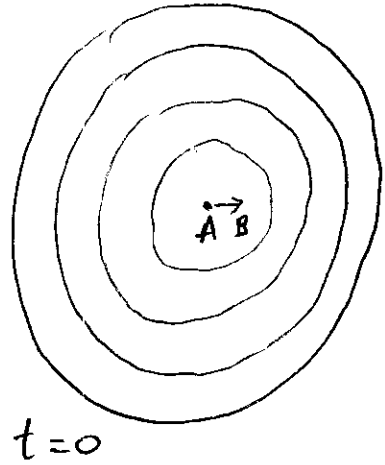
$\lambda_3 = \frac{2L}{3} \rightarrow$ in one L we have a wave length + a half wave length.

$\hookrightarrow \omega_3 = 3\omega_1$



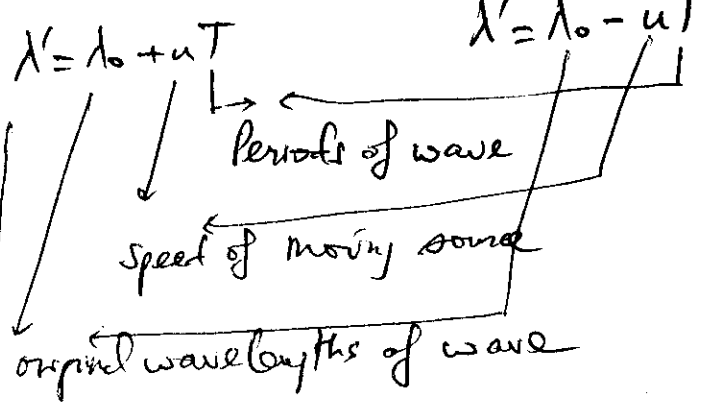
Doppler Effect: source of a wave is not static, it moves at constant speed u .

Wave takes some time to propagate \rightarrow wave emitted by moving source at previous time is still travelling which is caught up by a new wave emitted by the same source:



spread out $t=t_1$ wave fronts are further apart $\lambda' > \lambda_0$

piled up \rightarrow wave fronts are closer together $\lambda' < \lambda_0$



$$f' = \frac{f_0}{1 \pm \frac{u}{v}}$$

receding source

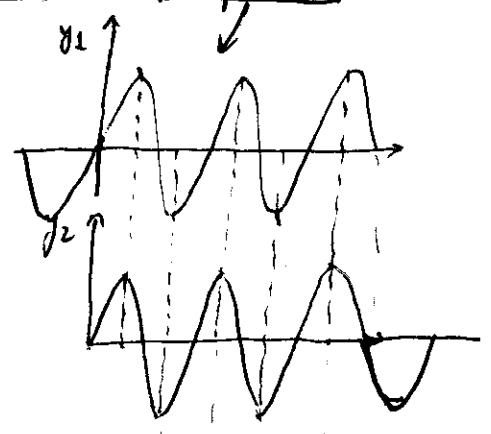
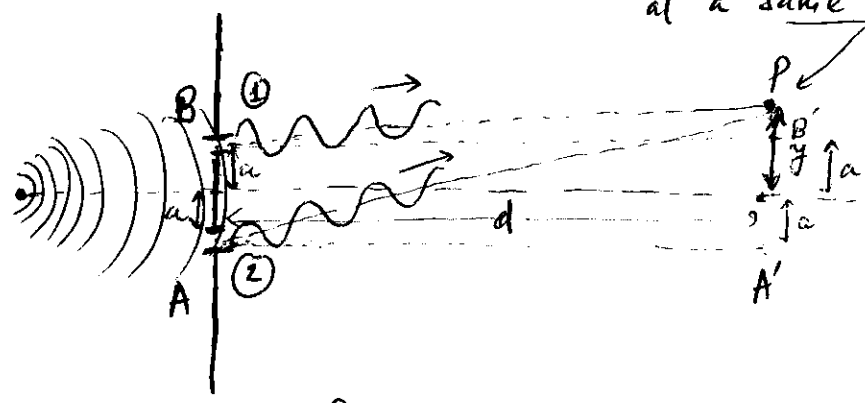
$$f' = \frac{f_0}{1 \mp \frac{u}{v}}$$

approaching source

f_0 : original frequency of wave ; v : wave propagation speed ; u : source speed.

Wave Motion (Cont.):
Constructive & Destructive Interference

Constructive Interference: Two identical waves ~~with~~ arriving at a same location **in phase**

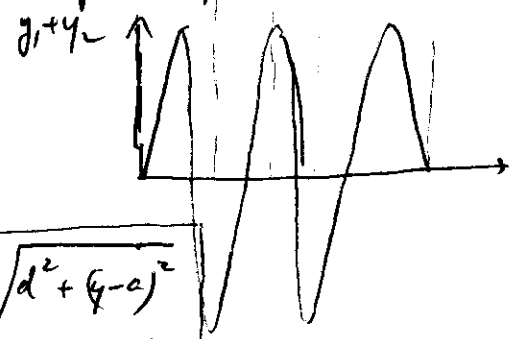


Two identical waves of wavelength λ travelling out from A & B towards P.

In phase:
Arrival.

$AP - BP = 0, \lambda, 2\lambda, 3\lambda, \dots$
 path difference = $n\lambda$ ($n = 0, 1, 2, \dots$)

AA'P : $AP = \sqrt{d^2 + (y+a)^2}$
 BB'P : $BP = \sqrt{d^2 + (y-a)^2}$



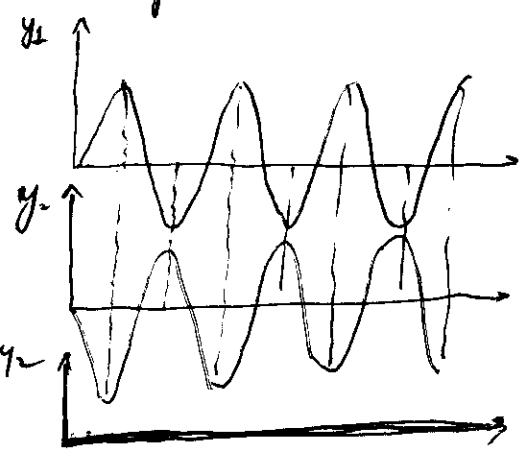
In-phase arrival: $AP - BP = \sqrt{d^2 + (y+a)^2} - \sqrt{d^2 + (y-a)^2} = n\lambda$ ($n = 0, 1, 2, 3, \dots$)
 (Constructive interference @ P)

P on midline b/w A & B.

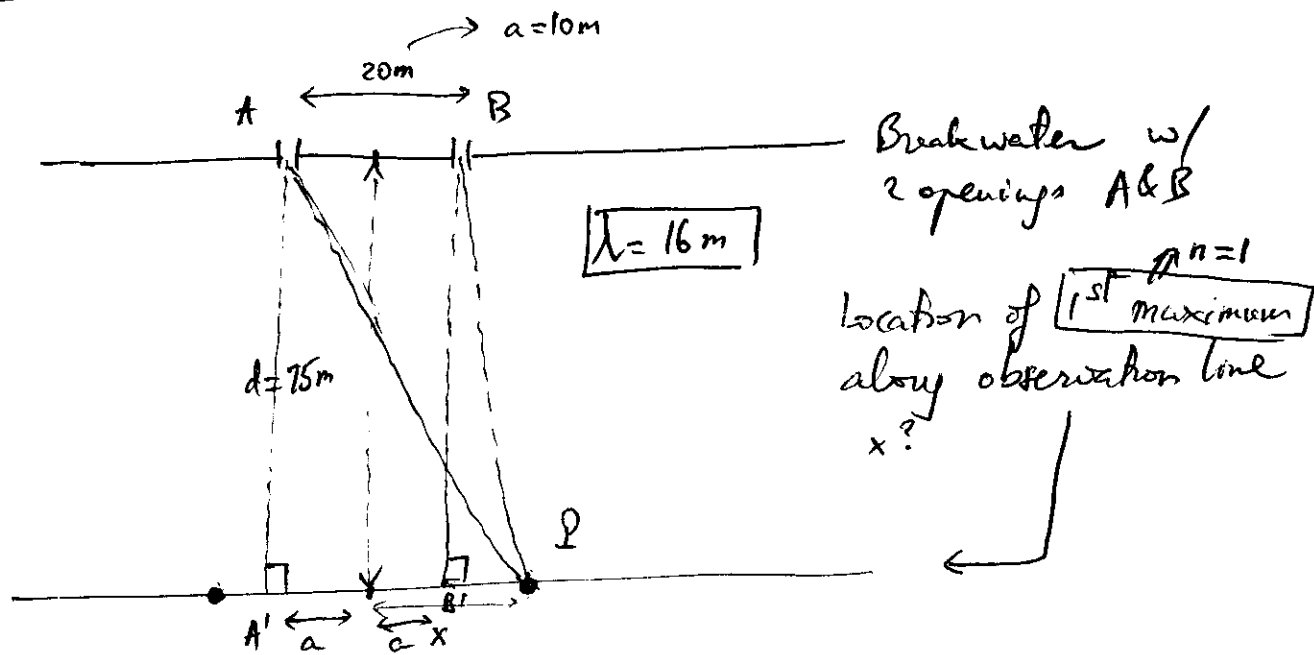
Destructive Interference: Two identical waves arriving at same location out of phase

Out of phase arrival:
 $AP - BP = (2n+1) \frac{\lambda}{2}$ ($n = 0, 1, 2, 3, \dots$)

$\sqrt{d^2 + (y+a)^2} - \sqrt{d^2 + (y-a)^2} = (2n+1) \frac{\lambda}{2}$
 odd number of half wave lengths
 $n = 0, 1, 2, 3, \dots$
 Destructive interference @ P.



Constructive & Destructive Interference in Water Wave:



1st max: $AP - BP = n\lambda = \lambda \quad (n=1)$

$$\sqrt{d^2 + (x+a)^2} - \sqrt{d^2 + (x-a)^2} = \lambda$$

$$\sqrt{75^2 + (x+10)^2} - \sqrt{75^2 + (x-10)^2} = 16$$

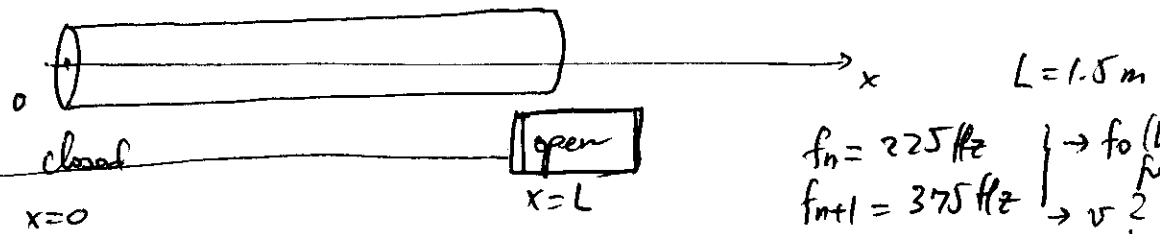
Solve for x: 1) Elevate both sides to power of 2
 2) Isolate cross term w/ square roots \rightarrow Repeat step 1.

$$75^2 + (x+10)^2 + 75^2 + (x-10)^2 - 2\sqrt{75^2 + (x+10)^2} - \sqrt{75^2 + (x-10)^2} = 16^2$$

$$[2 \cdot 75^2 + 2x^2 + 200 - 16^2]^2 = 4(75^2 + (x+10)^2)(75^2 + (x-10)^2)$$

Polynomial of order 4 $\rightarrow x = \pm 33\text{ m}$

14.72



$f_n = 225 \text{ Hz}$
 $f_{n+1} = 375 \text{ Hz}$

$\rightarrow f_0$ (lowest freq)
 $\rightarrow v$?

Pipe with one end closed \rightarrow sound waves

A closed end for sound waves is similar to a fixed end on a string \rightarrow Only certain wavelengths can stand in the pipe:

$$y_T = y_{inc} + y_{reflected}$$

$$= A \cos(kx - \omega t) - A \cos(kx + \omega t) = 2A \sin kx \sin \omega t$$

One end is closed $\rightarrow y_T(L, t) = 0 \rightarrow \sin kL = 0 \rightarrow kL = n\pi$

$$\frac{2\pi}{\lambda_n} L = n\pi \rightarrow \lambda_n = \frac{2L}{n} \quad (n = 1, 2, 3, \dots)$$

one end is ~~open~~ closed

\rightarrow Pipe with one end open: $\rightarrow y_T(L, t) = \text{max} \rightarrow \sin kL = 1$

$$kL = (2n+1) \frac{\pi}{2} \quad (n = 0, 1, 2, 3, \dots)$$

odd multiple of $\frac{\pi}{2}$

$$\frac{2\pi}{\lambda_n} L = (2n+1) \frac{\pi}{2} \rightarrow \lambda_n = \frac{4L}{2n+1} \quad (n = 0, 1, 2, 3, \dots)$$

one end is open.

Rewrite in terms of frequencies: $\left[v = \frac{\omega}{k} = \frac{2\pi f}{\frac{2\pi}{\lambda}} = f \cdot \lambda \right]$

$$f = \frac{v}{\lambda} \rightarrow \left[f_n = \frac{v}{\frac{4L}{2n+1}} = \frac{2n+1}{4L} v \quad (n = 0, 1, 2, 3, \dots) \right]$$

$$\frac{f_{n+1}}{f_n} = \frac{(2(n+1)+1)/4L \cdot v}{\frac{2n+1}{4L} v} = \frac{2n+3}{2n+1} = \frac{375}{225} = \frac{125 \times 3}{75 \times 3} = \frac{5}{3}$$

\rightarrow solve for $n \rightarrow \boxed{n=1}$

$f_1 = 225 \text{ Hz} \rightarrow f_2 = 375 \text{ Hz}$

$f_0?$

a) $\frac{f_1}{f_0} = \frac{\frac{3}{4} \lambda}{\frac{1}{4} \lambda} = 3 \rightarrow \boxed{f_0 = \frac{f_1}{3} = \frac{225}{3} = 75 \text{ Hz}}$

b) sound speed: v
 $f_1 = \frac{3}{4L} v \rightarrow \boxed{v = \frac{f_1 4L}{3} = \frac{225 \cdot 4 \cdot 1.5}{3} = 450 \frac{\text{m}}{\text{s}}}$

Ch 15 Fluid Motion

Fluid: Gas density ρ (rho) is variable \leftrightarrow compressible
 Liquid density ρ is constant \leftrightarrow incompressible

In a gas: molecules are far apart which can be brought closer together (compressing the gas)

Describe a fluid { Density $\rho = \frac{M \text{ (mass)}}{V \text{ (volume)}} = \frac{dm}{dV}$ (SI unit: $\frac{kg}{m^3}$)

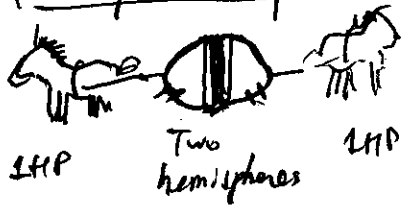
Pressure P : normal force per unit area $P = \frac{F}{A} \approx \frac{dF}{dA}$

(SI: $\frac{N}{m^2} = Pa$ or Pascal)
 other unit: Atm (Atmosphere)
 $1 \text{ Atm} = 1.013 \times 10^5 Pa$

Crushed barrel: experiment: a) open barrel was heated to a certain temperature so fewer gas molecules inside barrel will exert same pressure (atmospheric pressure) (most molecules were pushed out) lower ρ , same P
 b) barrel is closed tight (no further flow of gas molecules either way)
 c) ice were poured on barrel: lowering inside gas temperature: $\downarrow P$, same ρ

Outside barrel: atmospheric pressure
 Inside barrel: lower than atmospheric pressure
 \hookrightarrow Barrel is crushed by force of atmospheric pressure.

Hemisphere experiment:



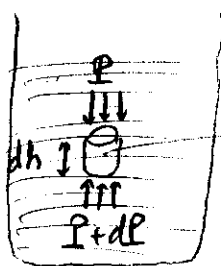
- a) Air from inside the hemisphere is pump out $\downarrow P$
- b) Higher P outside is keeping the hemisphere together

- Hydrostatic equilibrium \leftrightarrow water transportation
 - Conservation of mass
 - Conservation of energy
- } \rightarrow Bernoulli's principle
- gas \rightarrow air transportation
 - water \rightarrow hydro power plants.

Hydrostatic equilibrium:

fluid pressure compensating force of gravity
 \rightarrow floating objects in air (balloons) or in water (ships, boats, ...). Astronaut training in water

• The difference in pressure above (less) and below (more) on an element of fluid exactly cancels the weight of that element of fluid.



element of fluid of weight gdm and cross sectional area A & volume $dV = Adh$
 $\rightarrow dm = \rho dV$

Container filled with water on Earth
 P is higher at bottom $\leftrightarrow P$ is lower at top

$$A dP - gdm = 0$$

$$A dP - \rho g dV = 0$$

$$dP = \rho g dh$$

$\frac{dP}{dh} = \rho g$

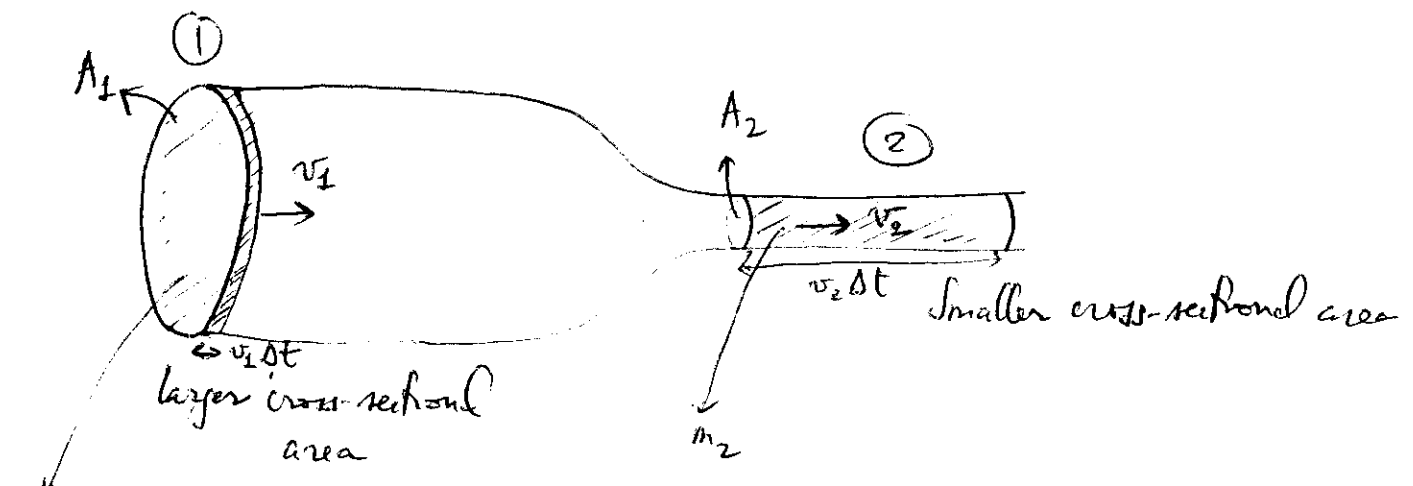
↓
 Buoyancy:
 $F_b = P \cdot A = \frac{\rho g h A}{\rho}$
 ↓
 buoyant force

Conservation of Mass:

→ Assuming no leaking

→ Consequence: $v \cdot A = \text{constant}$

\downarrow speed of fluid
 \downarrow cross-sectional area



Conservation of mass: $m_1 = m_2$

$$(\rho V)_1 = (\rho V)_2$$

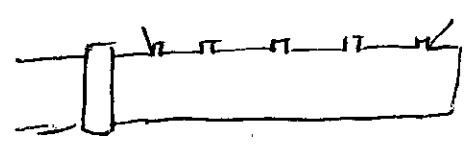
$$\rho v_1 \Delta t \cdot A_1 = \rho v_2 \Delta t \cdot A_2$$

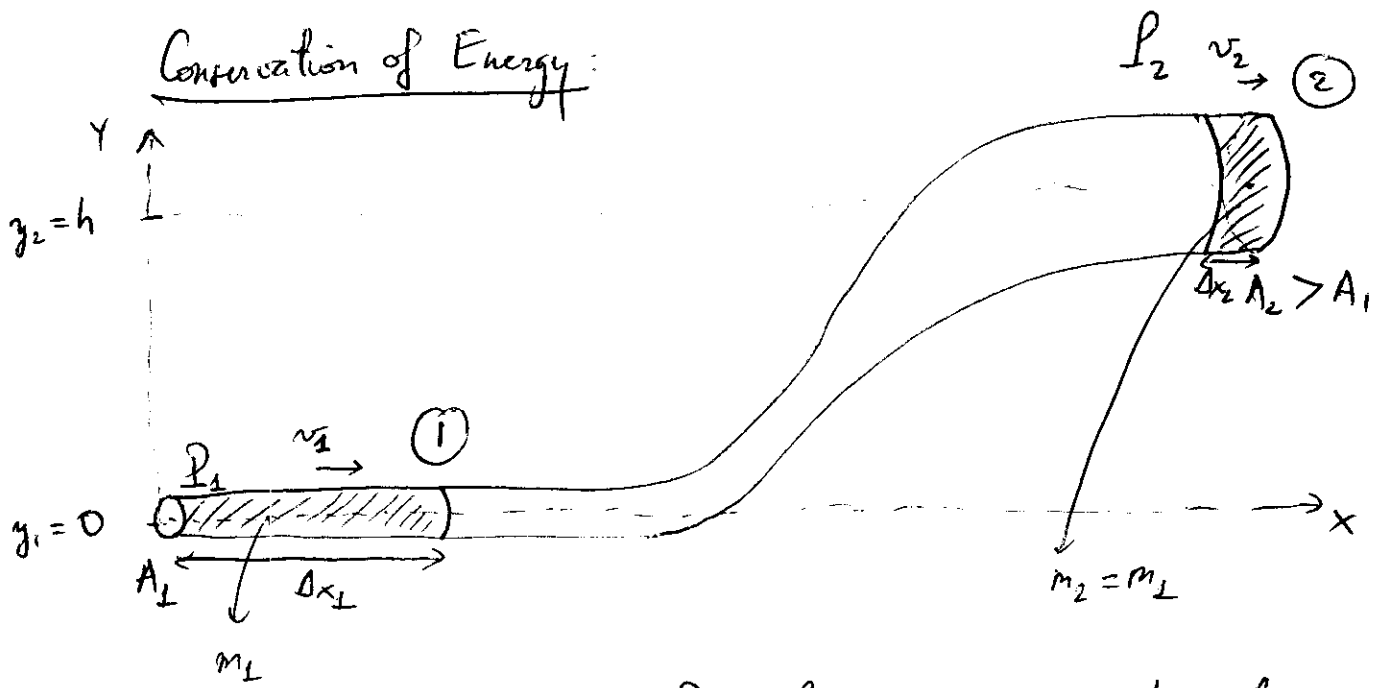
$$v_1 A_1 = v_2 A_2$$

$vA \text{ is constant}$

↳ larger cross-sect. area → lower speed
 Smaller " " " → higher speed.

Sprinkler:





This is possible only if $P_1 > P_2$: pressure has done work in moving column of liquid from ① to ②:

$$\Delta W = P_1 A_1 \Delta x_1 - P_2 A_2 \Delta x_2$$

Work done by pressure (F · Δx)

$$\left. \begin{aligned} & P_1 A_1 \Delta x_1 - P_2 A_2 \Delta x_2 = \\ & \frac{1}{2} m v_2^2 + m g y_2 - \frac{1}{2} m v_1^2 - m g y_1 \end{aligned} \right\}$$

↳ Is also $\Delta KE + \Delta P.E$

↳ Consequence: grouping terms at position ① & position ②:

$$\frac{1}{2} m v^2 + m g y + P A \Delta x = \text{constant}$$

Divide by volume: $V = A \Delta x$

$$\frac{1}{2} \left(\frac{m}{V} \right) v^2 + \left(\frac{m}{V} \right) g y + P \frac{A \Delta x}{A \Delta x} = \text{constant}$$

density

$$\frac{1}{2} \rho v^2 + \rho g y + P = \text{constant}$$

Bernoulli's equation.

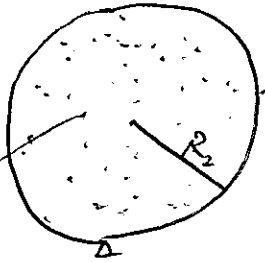
15.60

Helium balloon : buoyancy:

$$F_b = \rho h A \rightarrow$$

$$F_b = \rho V$$

Helium \rightarrow ①
Rubber \rightarrow ②



$m_2 = 0.85g$
(mass of rubber)
 $R_2 = 0.15m$

$$\rho_{He} = \rho_1 = 0.18 \frac{kg}{m^3}$$

How many paper clips (each has mass $m_3 = 1g$) can we hang before this balloon loses its buoyancy?

Buoyancy in this case: He has lower density than air

When N clips are attached: weigh by $\left\{ \begin{array}{l} \text{Rubber} \\ \text{He} \\ \text{Clips} \end{array} \right.$

$$(Nm_3 + m_2 + \underbrace{\rho_1 \frac{4}{3}\pi R_2^3}_{m_1}) \cdot g = F_b = \underbrace{\rho_{air} V_{balloon}}_{\text{mass of air displaced by He balloon}}$$

$$N = \frac{\rho_{air} \frac{4}{3}\pi R_2^3 - m_2 - \rho_{He} \frac{4}{3}\pi R_2^3}{m_3}$$

Data

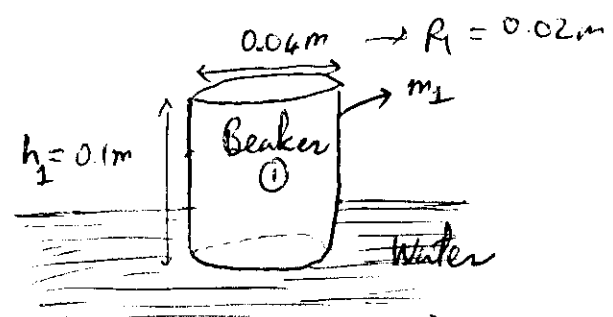
$$\rho_{air} = 1.2 \frac{kg}{m^3}$$

$$N = \frac{(1.2 - 0.18) \frac{4}{3}\pi (0.15)^3 - 0.85 \cdot 10^{-3}}{10^{-3}} = 13.6 \text{ clips}$$

↓
13 clips before balloon loses buoyancy.

15.50

Similar to 15.60 but now in a liquid (but also a fluid!)



Beaker floating in water with $\frac{1}{3}$ submerged:

weight of beaker is held by a buoyant force F_b

$$F_b = \rho_{\text{water}} V g \quad \text{where } V \text{ vol. of water displaced by beaker}$$

(in 15.60: V here was the volume of the ^{entire} balloon since whole balloon was submerged in air)

$$\text{in 15.50: } V = \frac{1}{3} V_1 = \frac{1}{3} \pi R_1^2 h_1$$

$$\rightarrow F_b = \rho_{\text{water}} \frac{1}{3} \pi R_1^2 h_1 g = m_1 g \rightarrow \text{here } m_1 \text{ is obtained.}$$

How many rocks (each of mass $m_2 = 15\text{g}$) can be placed in beaker before it sinks \leftrightarrow if we submerge another $\frac{2}{3}$ of beaker can hold N rocks: additional buoyant force is \tilde{F}_b

$$\tilde{F}_b = \rho_{\text{water}} \frac{2}{3} \pi R_1^2 h_1 g = N m_2 g$$

$$N = \frac{\rho_{\text{water}} \frac{2}{3} \pi R_1^2 h_1 g}{m_2 g}$$

$$= \frac{1000 \frac{2}{3} \pi \cdot 0.02^2 \cdot 0.1}{15 \cdot 10^{-3}}$$


$$N = 5.59$$

\downarrow
5 rocks before beaker sinks.

Data:

$$\rho_{\text{water}} = \frac{1000 \text{ kg}}{\text{m}^3}$$

15.22

4300kg elephant on one foot approximated as a circle  R = 0.3m

$$P = ? = \frac{F}{\text{Area}} = \frac{m \cdot g}{\pi R^2} = \frac{4300 \cdot 9.81}{\pi \cdot 0.3^2} = 596 \frac{\text{kPa}}{10^3}$$

(1Atm = normal atmospheric pressure = 101300 Pa
 ≈ 101 kPa)

14.45

Exam 3 Review of Topics (Ch 10-15)

Rotational Motion (Ch. 10)

Torque: select pivot point $\vec{\tau} = \vec{r} \times \vec{F} = rF \sin \theta \hat{e}$
 position vector \vec{r} : from pivot point to force application point
 θ = angle b/w \vec{r} & \vec{F}
 \hat{e} : perpendicular to both \vec{r} & \vec{F} , given by RHR

Analogy's of 2nd Newton's Law

$$\vec{\tau} = I \cdot \vec{\alpha}$$

I : moment of inertia
 $\vec{\alpha}$: vector angular acceleration.

Angular Momentum (Ch. 11)

$$\vec{L} = \vec{r} \times \vec{p} \quad (\vec{p} = \text{linear momentum } \vec{p} \equiv m\vec{v})$$

$$= r p \sin \alpha \hat{L}$$

\vec{r} : position vector: from pivot to position of object
 α : angle b/w \vec{r} & \vec{p}
 \hat{L} : perpendicular to both \vec{r} & \vec{p} , given by RHR

Round object: $L = I\omega$ (I : m. of inertia; ω : angular speed)

Conservation of Angular momentum:

$$\vec{\tau}_{\text{net external}} = \frac{d\vec{L}}{dt}; \quad \vec{\tau}_{\text{net external}} = 0 \rightarrow \boxed{L_i = L_f}$$

Static Equilibrium (Ch. 12)

Conditions: $\sum \vec{F}_i = 0$
 $\sum \vec{\tau}_i = 0 \leftrightarrow$ select pivot so force with least inflection offer no torque.
 Use geometry to find angles b/w \vec{r} & \vec{F} for different torques. Torque directions by RHR.

Oscillatory Motion (Ch. 13)

SHM {

- Pendulum (Gravitational): $\frac{d^2\theta}{dt^2} = -\frac{g}{L}\theta \rightarrow \theta(t) = \theta_m \cos \omega t$
 $\omega = \sqrt{\frac{g}{L}}$ (small angle approx.)
- Torsional Pendulum: $\frac{d^2\theta}{dt^2} = -\frac{K}{I}\theta \rightarrow \theta(t) = \theta_m \cos \omega t$
 $\omega = \sqrt{\frac{K}{I}}$
- Spring & bob: $\frac{d^2x}{dt^2} = -\frac{k}{m}x \rightarrow x(t) = x_m \cos \omega t$
 $\omega = \sqrt{\frac{k}{m}}$

Total energy is constant: $\frac{1}{2}mv^2 + \frac{1}{2}kx^2 = \frac{1}{2}kx_m^2$ (constant).

Wave Motion (Ch 14)

Types {

- Longitudinal: osc. in same direction as propagation.
- Transverse**: osc. perpendicular to direction of propagation.

Wave in a string: {

- Math description: $y(x,t) = A \sin(kx - \omega t)$
- propagation in +x $\leftrightarrow A \sin(kx - \omega t)$
- propagation in -x $\leftrightarrow A \sin(kx + \omega t)$
- $v = \sqrt{\frac{T}{\mu}}$ { v : wave speed
 T : tension in string
 μ : linear mass density of string.

Properties:

Superposition

Beat phenomenon: amplitude modulation $\propto \cos\left(\frac{\omega_1 - \omega_2}{2}t\right)$ or beats (low freq.)

Intensity $\propto \cos^2 \rightarrow$ frequency is $\omega_1 - \omega_2$

Standing waves: Inc + reflected $2A \sin kx \sin \omega t$

- pipe one ^{closed} end: $x=L$ $f_n = \frac{n}{2L}v$ ($n=1,2,3,\dots$)
- pipe one ^{open} end: $x=L$ $f_n = \frac{2n+1}{4L}v$ ($n=0,1,2,3,\dots$)

Interference {

- Constructive: in phase arrival: $AP - BP = \sqrt{d^2 + (y+a)^2} - \sqrt{d^2 + (y-a)^2} = n\lambda$ ($n=0,1,2,3,\dots$)
- Destructive: out of phase arrival: $AP - BP = (2n+1)\frac{\lambda}{2}$ ($n=0,1,2,3,\dots$)

Doppler effect: source is moving @ speed u :

$$f' = \frac{f_0}{1 \pm \frac{u}{v}}$$

receding source \leftarrow \rightarrow approaching source

Fluid Motion (Ch 15)

- Hydrostatic Equilibrium : $\frac{dP}{dh} = \rho g \rightarrow$ Buoyancy $= F_b = \rho g V$

Buoyant force on a He balloon of volume V is given by the amount of air displaced by the balloon: $\rho_{air} V_{balloon}$

- Conservation of mass : $vA = \text{constant}$
 ↳ fluid speed ↳ cross sectional area of pipe

- Conservation of energy : $\frac{1}{2} \rho v^2 + \rho gh + P = \text{constant}$
 Bernoulli's equation.

$P = \frac{F}{A}$; SI : Pa
 1 Atm = 101300 Pa

14.45

Doppler Effect:

$$f' = \frac{f_0}{1 + \frac{u}{v}} \rightarrow \frac{f'}{f_0} = \frac{1}{1 + \frac{u}{v}}$$

receding source : $f' < f_0$

$$\boxed{f = \frac{v}{\lambda}} \rightarrow \lambda = \frac{v}{f}$$

$$\frac{\lambda}{T} = v$$

$$\lambda \propto \frac{1}{f} \rightarrow \frac{\lambda'}{\lambda_0} = 1 + \frac{u}{v}$$

Distant galaxy speed.

→ wave → light speed.

$$\frac{708 \text{ nm}}{656 \text{ nm}} = 1 + \frac{u}{c} \rightarrow u = \left(\frac{708}{656} - 1 \right) \cdot 3 \times 10^8 = 2.38 \times 10^7 \frac{\text{m}}{\text{s}}$$